DERIVATION OF IMPLICIT RATIONAL RUNGE-KUTTA SCHEMES FOR APPROXIMATION OF SECOND ORDER ORDINARY DIFFERENTIAL EQUATIONS

 \mathbf{BY}

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CERTIFICATION

This is to certify that this thesis "Derivation of Implicit Rational Runge – Kutta Schemes for the Approximation of Second Order Ordinary Differential Equations" has been duly presented by SANI, Usman Abdullahi (M.Sc/MA/06/0118) of the department of Mathematics of Modibbo Adama University of Technology, Yola, and been approved by the following having met the stipulated requirement.

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DEDICATION

This work is dedicated to my parent, Alh. Abdullahi Sani and Haj. Zainabu Abu.

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Alhandulillah! To the only wonderful King, who gave me the privilege and life from the beginning of this work up to this stage. Honour and glory be the Almighty Allah.

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ABSTRACT

The development of 3 – stage Implicit Rational Runge – Kutta methods are considered using Taylor and Binomial series expansion for the direct solution of general second order initial value problems of ordinary differential equations with constant step length. The basic properties of the developed method were investigated and found to be consistent and convergent. The efficiency of the method were tested on some numerical examples and found to give better approximations than the existing methods.

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CHAPTER ONE

INTRODUCTION

1.0 Background of the Study

A numerical method is the study of approximate techniques for solving mathematical problems taking into account the extent of possible errors. Many numerical schemes have been developed for solution of ordinary differential equations which posses no analytic solution. In recent years much research has been carried out to develop derivative free Runge-Kutta methods.

Conventionally, higher order ordinary differential equations are solved by method of reduction to system of first order ordinary differential equations and methods of solving the resultant first order ordinary differential equations can be adopted. The approach has reported to increase the dimension of the problem and therefore result in more computational burden. Awoyemi (2001), Adesanya *et al.* (2009) and Anake *et al.* (2012). Some of the direct methods of solving second order ordinary differential equations are: Nystrom type, self stating Runge – Kutta method which involve several functions evaluation per step and the linear multi-step methods, particularly the Implicit methods which have better stability than the explicit methods and require fewer functions evaluation per step, (Jator (2010), Agular and Ramos (2006)).

1.1 Statement of the Problem

Most of the existing Runge – Kutta methods were developed to solve first order initial value problems or special second order ordinary differential equations. This thesis considered the direct solution of general second order initial value problems.

1.2 Aim and Objectives of the Study

The aim of this study is to develop a 3 – Stage Implicit Rational Runge – Kutta method for the direct solution of general second order initial value problems, while the objectives are:

- i. Use Taylor and Binomial series expansion of the s-stage Rational Runge-Kutta method to derive an implicit rational Runge – Kutta schemes for second order initial value problems.
- ii. Analyze the consistency, convergence and the error analysis of the schemes.
- iii. Verify the efficiency of the derived scheme using numerical examples.

1.3 Significance of the Study

The significance of this study is to derive a better scheme in terms of approximation and time of execution for the direct solution of general second order ordinary differential equations.

1.4 Scope of the study

The method developed in this thesis can only solve second order initial value problems.

1.5 Definition of terms

Derivative:

The rate of change of one variable with respect to another is called derivative.

Differential equation (D.E):

An equation involving some of the derivative of functions is called differential equation.

Type of differential equations:

An equation that contains one independent variable and its derivative with respect to that independent variable is called Ordinary Differential Equation (ODE). E.g

$$\frac{d^2y}{dx^2} + 5\frac{dy}{dx} + 3y = \sin x$$

while on the other hand, if the equation contains more than one independent variable and their derivatives with respect to each independent variables is called Partial Differential Equations (PDE). E.g

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

Order of Differential equation:

The order of D.E is the order of the highest derivative of the unknown functions that appear in the equation. Thus, the equation

$$\frac{d^2y}{dx^2} - 6\frac{dy}{dx} + 2y = 0$$
 is of order 2.

Linear and Non - Linear Differential Equation

Consider the general form of ODE

$$a_n y^n + a_{n-1} y^{n-1} + a_{n-2} y^{n-2} + \dots + a_0 y = f(t)$$

If the coefficients a_n , a_{n-1} , ..., a_{0} , are either constants or functions of t, then the ODE is said to be linear, otherwise is said to be non-linear.

Consistency

The Runge-Kutta method is said to be consistent with the initial value condition if

$$\emptyset(x, y, 0) \equiv f(x, y) \tag{1.10}$$

then,

$$y(x+h) - y(x) - h\emptyset(x, y(x), h) = y(x) + hy'(x) + O(h^2) - y(x) - h\emptyset(x, y(x), 0)$$
$$= hy'(x_n) - hy'(x_n) + O(h^2) = O(h^2)$$

Since $y'(x) = f(x, y(x)) = \emptyset(x, y(x), 0)$ by the above then a consistent method has order at least one, Lambert (1973).

Order

A method is said to have order p if p is the largest integer for which

$$y(x + h) - y(x) - h\emptyset(x, y(x), h) = O(h^{p+1})$$

holds, the Taylor algorithm of order p is

$$y(x + h) = y(x) + hy'(x) + \frac{h^2}{2!}y''(x) + \dots + \frac{h^p}{p!}y^p(x)$$

Stability

A numerical method is stable if small change in the initial condition produces a corresponding small change in the subsequent approximations.

A – Stability

A numerical method is said to be "A-Stable" if its region of absolute stability contains the whole of the left hand half plane i.e. $Reh\lambda < 0$.

Convergent

A numerical method is said to be convergent if the numerical solution approaches the exact solution as the step size goes to zero. i.e.

$$\lim_{h \to 0} [y(x_{n+1}) - y_{n+1}] = 0$$

Stiff Equation

The initial value problem y' = f(x, y) is considered stiff when its Jacobian matrix has both eigenvalues with very large negative real part and eigenvalues with very small negative real part which constrain the step length to be extremely small to maintain stability requirements in Numerical methods.

Truncation Error

Truncation error: an error introduced as a result of ignoring some of the higher terms of the power series expansion during the development of the new formula. Mathematically written as

$$T_{n+1} = |y(x_{n+1}) - y(x_n) - h\phi(x_n, y_n, h)|$$

Explicit Scheme

In an explicit numerical scheme, the calculation of the solution at a given step or stage does not depend on the value of the solution at that step or on a later step or stage.

Implicit Scheme

In an implicit numerical scheme, the calculation of the solution at a given step or stage does depend on the value of the solution at that step or on a later step or stage. Such methods are usually more expensive than explicit schemes but are better for handling stiff ODEs.

Embedded Runge-Kutta methods

Two Runge-Kutta methods that share the same stages. The differences between the solutions give an estimate of the local truncation error.

CHAPTER TWO

LITERATURE REVIEW

2.0 Introduction

In numerical analysis, Runge-Kutta schemes are important family of implicit and explicit iterative methods for approximation of solution of ordinary

differential equations. The method was developed around 1900 by German Mathematician C. Runge and M. W. Kutta.

First order initial value problem is given in the form:

$$y' = f(x, y), y(x_0) = y_0$$
 (2.1)

The numerical solution to (2.1) is

$$y_{n+1} = y_n + h\phi(x_n, y_n, h)$$
 (2.2)

where

$$\phi(x, y, h) = \sum_{i=1}^{3} c_i k_i$$

$$k_1 = f(x, y),$$

$$k_r = f\left(x + ha_i, y + h\sum_{j=1}^{r} b_{ij} k_j\right), \quad r = 1(1)s$$
(2.3)

with constraints

$$a_i = \sum_{j=1}^i b_{ij}, \quad i = 1(1)s$$

The derivative of suitable parameters a_{ij} , b_i and c_i of higher order term involves a large amount of tedious algebraic manipulations and functions evaluations which is both time consuming and error prone, Julyan and Oreste (1992). The derivation of the second order, third order, otherwise known as the Heun's method and the popular fourth order Runge – Kutta methods is extensively discussed by Lambert (1973), Butcher (1987). According to Julyan and Oreste (1992) the minimum number of stages necessary for an explicit method to attain order p is still an open problem. Therefore so many new

schemes and approximation formula have been derived this includes the work of Ababneh *et al.* (2009a), Ababneh *et al.* (2009b) Faranak and Ismail (2010).

The Runge-Kutta method is not restricted to solving first order differential equation of the form (2.1). We can do the same by splitting higher order differential equations into an equivalent system of first order equations.

In general, the second order equation can be reduced to an equivalent system of first order of twice the dimension and solved using the standard Runge-Kutta method. However, it is more efficient if the equation can be solved directly using Runge-Kutta-Nystrom method of the special second order differential equations which according to Jain (1984) is given in the form:

$$y'' = f(x, y),$$
 $y(x_0) = y_0, y'(x_0) = y_0$ (2.4)

with numerical solution

$$y_{n+1} = y_n + hy'_n + h^2 \sum_{i=1}^{s} b_i k_i$$
 (2.5a)

and

$$y'_{n+1} = y'_n + h \sum_{i=1}^{s} b'_i k_i$$
 (2.5b)

where

$$k_i = f\left(x_n + c_i h, y_n + c_i h y'_n + h^2 \sum_{j=1}^i a_{ij} k_j\right), \quad i = 1, 2, ..., s$$

The derivative of suitable parameters a_{ij} , b_i and c_i requires extremely lengthy algebraic manipulations, except for small values of s. Such method was discussed in Sharp and Fine (1992) and Dormand $et\ al\ (1987)$.

Runge-Kutta-Nystrom methods are direct extension of Runge-Kutta method to second order differential equations in (2.4) so given by the Finnish Mathematician E. J Nystrom. The best known of these method uses the formula in (2.5) where h = 0, 1, 2, ..., N - 1 (N, the number of step), Nystrom (1925).

2.1: Runge – Kutta Nystrom Methods

Recently, significant advances have been made in the development of general theory of Runge-Kutta-Nystrom methods and in the derivation, error estimation and control, Dormand (1996) which includes many notable ones.

Sharp and Fine (1992) generated some Nystrom pairs for the general second order initial value problem. Dormand *et al.* (1987) obtained some families of Runge-Kutta-Nystrom formula. Fudziah (2003) derived the embedded pairs of Runge-Kutta-Nystrom method which are diagonally implicit and all the diagonal elements are equal. Fudziah (2009) also derived a Sixth Order Singly Diagonally Implicit Runge-Kutta Nystrom Method with Explicit First Stage for Solving Second Order Ordinary Differential Equations. Senu *et al.* (2011) derived a Singly Diagonally Implicit Runge-Kutta - Nystrom Method for Solving Oscillatory Problems. Also recently Okunuga *et al.* (2012) discussed the general techniques for solving equation of the form (2.4) directly without first reducing it to systems of first order ODEs. The above authors observed that the direct solution of second order equations is of greater advantage over reduction to systems of first order equations to increase efficiency and reduced storage requirement.

It should be noted that the methods considered above were Runge-Kutta-Nystrom methods of the special cases for second order ordinary differential equation of the form (2.4). Much attention have not been given to Runge – Kutta method for the solution of general second order ODEs; bearing in mind

that many physical systems are modeled using the general second order ODEs, Lambert (1973).

2.2: Runge-Kutta Method for Second Order Ordinary Differential Equations.

The general s – stage Runge-Kutta scheme for general second order initial value problems of ordinary differential equations of the form

$$y'' = f(x, y, y'),$$
 $y(x_0) = y_0, y'(x_0) = y_0$ (2.2.1)

is defined by Jain (1984) to be

$$y_{n+1} = y_n + hy_n' + \sum_{r=1}^{s} c_r k_r$$
 (2.2.2)

and

$$y'_{n+1} = y'_n + \frac{1}{h} \sum_{r=1}^{s} c'_r k_r$$
 (2.2.3)

where

$$K_{r} = \frac{h^{2}}{2} f\left(x_{n} + c_{i}h, y_{n} + hc_{i}y_{n}' + \sum_{j=1}^{r} a_{ij}k_{j}, y_{n}' + \frac{1}{h}\sum_{j=1}^{r} b_{ij}k_{j}\right), \quad i = 1(1)s \quad (2.2.4)$$

with

$$c_i = \sum_{j=1}^i a_{ij} = \frac{1}{2} \sum_{j=1}^i b_{ij}, \qquad i = 1(1)r$$
 (2.2.5)

where c_i , a_{ij} , b_{ij} , c_r , c_r' are constants to be determined.

The fourth order Runge-Kutta method for the solution of (2.2.1) is:

$$y_{n+1} = \frac{1}{3}(k_1 + k_2 + k_3)$$
 , and

$$y'_{n+1} = \frac{1}{3h} (k_1 + 2k_2 + 2k_3 + k_4)$$
 (2.2.6)

where

$$k_{1} = h^{2} f(x, y, y')$$

$$k_{2} = \frac{h^{2}}{2} f\left(x + \frac{1}{2}h, y + \frac{h}{2}y' + \frac{1}{2}k_{1}, y' + \frac{1}{h}k_{1}\right)$$

$$k_{3} = \frac{h^{2}}{2} f\left(x + \frac{1}{2}h, y + \frac{h}{2}y' + \frac{1}{4}k_{2}, y' + \frac{1}{h}k_{2}\right)$$

$$k_{4} = \frac{h^{2}}{2} f\left(x + h, y + hy' + \frac{1}{4}k_{3}, y' + \frac{2}{h}k_{3}\right)$$

2.3: Rational Runge – Kutta Methods

Hong (1982) first proposed the use of rational function of Runge – Kutta method then Okunbor (1987) investigated the use of the rational function of the Runge-Kutta scheme:

$$y_{n+1} = \frac{y_n + h \sum_{i=1}^r w_i k_i}{1 + h y_n \sum_{i=1}^r v_i H_i}$$
 (2.3.1)

where

$$k_i = f(x_n + c_i h, y_n + h \sum_{i=1}^r a_{i-1j} k_j), i = 1(1)r$$
 (2.3.2)

and

$$H_i = g(x_n + d_i h, z_n + h \sum_{i=1}^r b_{i-1j} H_j), \ i = 1(1)r$$
(2.3.3)

in which

$$g(x_n, z_n) = -z_n^2 f(x_n, y_n)$$
 and $z_n = \frac{1}{y_n}$

where $c_{i,a_{ij}}$, b_{ij} , d_{i} are arbitrary constants to be determined.

The scheme is said to be

- i. Explicit if $b_{ij} = 0, j \ge i$
- ii. Semi-implicit if $b_{ij} = 0, j > i$
- iii. Implicit if $b_{ij} \neq 0$ for at least one j > i, h is the step size and the constraint

$$d_i = \sum_{i=1}^i b_{ij} \,, \tag{2.3.4}$$

is imposed to ensure consistency of the method.

The explicit scheme can only solve non-stiff and mildly stiff equations effectively. They performed poorly on stiff equations.

In view of these inadequacies of the explicit schemes and the superior region of absolute stability associated with implicit schemes, Ademuluyi and Babatola (2000) redefined the parameters; a_{ij} , b_{ij} , v_i , w_i , d_i , and c_{irs} so that the resulting method to be implicit and generates also the parameters so that the resulting numerical approximation method shall be A-stable and will have low bound for local truncation error. The scheme was used to solve some stiff initial value problems.

Odekunle (2001) generated some semi-implicit rational Runge-Kutta scheme which is also A-stable and can solve mildly and slightly stiff problems and reduces problems in solving coupled simultaneous non linear systems when compared with implicit schemes. Odekunle *et al.* (2004) developed a method that takes into consideration the point of singularities which other methods mentioned above cannot solve.

Since then many new rational Runge –Kutta schemes have been developed for the solution of first order initial value problems and found out to give better estimates. Among these authors are:

Bolarinwa (2005), who derived some class of semi-implicit rational Runge–Kutta schemes for solving ordinary differential equations with derivatives discontinuities, Babatola *et al.* (2007) developed a One-Stage Implicit rational Runge-Kutta schemes for the solution of discontinuous initial value problems and found out to be consistent, A-Stable and convergent.

Bolarinwa *et al.* (2012) developed a Two-Stage Semi-Implicit Rational Runge – Kutta Scheme for solving first order ordinary differential equations. The scheme is absolutely stable, consistent and convergent and was used to approximate a variety of first order differential equations. However, the methods are presently receiving more attention as efficient schemes for the solutions of various types of first order initial value problems.

The rational function of Runge-Kutta method is still open for more approximation schemes, estimation and control.

Soomiyol (2011) investigated the use of rational form of (2.2.2) and (2.2.3), defined by

$$y_{n+1} = \frac{y_n + hy_n' + \sum_{r=1}^{s} w_r K_r}{1 + y_n' \sum_{r=1}^{s} v_r H_r}$$
(2.3.5)

and

$$y'_{n+1} = \frac{y_n + \frac{1}{h} \sum_{r=1}^{s} w'_r K_r}{1 + \frac{1}{h} y'_n \sum_{r=1}^{s} v'_r H_r}$$
(2.3.6)

where

$$K_{r} = \frac{h^{2}}{2} f\left(x_{n} + c_{i}h, y_{n} + hc_{i}y_{n}' + \sum_{j=1}^{r} a_{ij}K_{j}, y_{n}' + \frac{1}{h}\sum_{j=1}^{r} b_{ij}K_{j}\right), i = 1(1)s$$
(2.3.7)

and

$$H_{r} = \frac{h^{2}}{2} g \left(x_{n} + d_{i}h, z_{n} + hd_{i}z'_{n} + \sum_{j=1}^{r} \alpha_{ij}H_{ij}, z'_{n} + \frac{1}{h}\sum_{j=1}^{r} \beta_{ij}H_{j} \right), i = 1(1)s$$
(2.3.8)

with constraints

$$c_i = \sum_{j=1}^i a_{ij} = \frac{1}{2} \sum_{j=1}^i b_{ij}, \quad i = 1(1)r$$

$$d_i = \sum_{j=1}^i \alpha_{ij} = \frac{1}{2} \sum_{j=1}^i \beta_{ij}, \qquad i = 1(1)r$$

in which

$$g(x_n, z_n, z'_n) = -z_n^2 f(x_n, y_n, y'_n)$$
 and $z_n = \frac{1}{v_n}$

The above scheme is said to be

- i. Explicit if a_{ij} , $b_{ij} = 0$, $j \ge i$
- ii. Semi-Implicit if a_{ij} , $b_{ij} = 0$, j > i
- iii. Implicit if a_{ij} , $b_{ij} \neq 0$, for at leasst one j > i

For direct solution of general second order ordinary differential equation and derived the Explicit Rational Runge-Kutta schemes. The scheme is found to be efficient, and was used to solve second order initial value problems. The methods mentioned above are explicit, hence are unsuitable for stiff systems; then there is the need to develop a method that is implicit in nature which have

better stability properties and able to solve stiff problems. In this thesis, we addressed the problem by developing implicit rational Runge-Kutta schemes.

CHAPTER THREE

METHODOLOGY

3.0 RATIONAL RUNGE-KUTTA SCHEME FOR SECOND ORDER ORDINARY DIFFERENTIAL EQUATIONS

In this chapter, the derivation of one stage, two stage and three stage implicit rational Runge-Kutta schemes is considered using the techniques of rational form of Runge-Kutta methods proposed by Okunbor (1987).

The *s*-stage rational Runge-Kutta method as proposed by Soomiyol (2011) is given as follows:

$$y_{n+1} = \frac{y_n + hy_n' + \sum_{r=1}^{s} w_r K_r}{1 + y_n' \sum_{r=1}^{s} v_r H_r}$$
(3.1)

$$y'_{n+1} = \frac{y_n + \frac{1}{h} \sum_{r=1}^{s} w'_r K_r}{1 + \frac{1}{h} y'_n \sum_{r=1}^{s} v'_r H_r}$$
(3.2)

where

$$K_{r} = \frac{h^{2}}{2} f\left(x_{n} + c_{i}h, y_{n} + hc_{i}y_{n}' + \sum_{j=1}^{r} a_{ij}K_{j}, y_{n}' + \frac{1}{h}\sum_{j=1}^{r} b_{ij}K_{j}\right), i = 1(1)s , \quad (3.3)$$

$$H_{r} = \frac{h^{2}}{2} g \left(x_{n} + d_{i}h, z_{n} + hd_{i}z'_{n} + \sum_{j=1}^{r} \alpha_{ij}H_{j}, z'_{n} + \frac{1}{h} \sum_{j=1}^{r} \beta_{ij}H_{j} \right), i = 1(1)s$$
(3.4)

with constraints

$$c_{i} = \sum_{j=1}^{i} a_{ij} = \frac{1}{2} \sum_{j=1}^{i} b_{ij}, \quad i = 1(1)r$$
(3.5)

$$d_{i} = \sum_{j=1}^{i} \alpha_{ij} = \frac{1}{2} \sum_{j=1}^{i} \beta_{ij}, \quad i = 1(1)r$$
(3.6)

in which

$$g(x_n, z_n, z'_n) = -z_n^2 f(x_n, y_n, y'_n)$$
 and $z_n = \frac{1}{y_n}$ (3.7)

The constraint equations (3.5) and (3.6) are to ensure consistency of the method, h is the step size and the parameters a_{ij} , b_{ij} , c_i , $d_{i,\alpha_{ij}}$, β_{ij} are constants called the parameters of the method.

In this research work, we shall consider the implicit scheme where all the a_{ij} and $b_{ij} \neq 0$ for at least one j > i.

3.1 DERIVATION OF ONE STAGE IMPLICIT SCHEME

From equations (3.1), (3.2), (3.3) and (3.4), setting s = 1, we have

$$y_{n+1} = \frac{y_n + hy_n' + w_1 K_1}{1 + y_n' v_1 H_1}$$
(3.1.1)

and

$$y'_{n+1} = \frac{y_n + \frac{1}{h}w'_1K_1}{1 + \frac{1}{h}y'_nv'_1H_1}$$
(3.1.2)

where

$$k_1 = \frac{h^2}{2} f\left(x_n + c_1 h, y_n + h c_1 y'_n + a_{11} K_1, y'_n + \frac{1}{h} b_{11} K_1\right), i = 1(1)s \quad (3.1.3a)$$

and

$$H_1 = \frac{h^2}{2}g\left(x_n + d_1h, z_n + hd_1z'_n + \alpha_{11}H_1, z'_n + \frac{1}{h}\beta_{11}H_1\right), i = 1(1)s \quad (3.1.3b)$$

with constraints

$$c_1 = a_{11} = \frac{1}{2}b_{11}$$
 and $d_1 = \alpha_{11} = \frac{1}{2}\beta_{11}$ (3.1.4)

where c_1 , a_{11} , b_{11} , d_1 , α_{11} , β_{11} , w_1 , w_1' , w_1' , v_1 and v_1' are all constants to be determined.

Equation (3.1.1) can be written as

$$y_{n+1} = (y_n + hy_n' + w_1k_1)(1 + y_nv_1H_1)^{-1}$$
(3.1.5)

Expanding the bracket and neglecting 2nd and higher orders gives

$$y_{n+1} = (y_n + hy_n' + w_1k_1)(1 - y_nv_1H_1)$$
(3.1.6)

Expanding (3.1.6) and re-arranging, gives

$$y_{n+1} = y_n + hy'_n - (y_n^2 v_1 + hy_n y'_n v_1) H_1 + (w_1 - y_n v_1 H_1 w_1) K_1$$
(3.1.7)

Equation (3.1.2) can be written as

$$y'_{n+1} = \left(y'_n + \frac{1}{h}w'_1K_1\right)\left(1 + \frac{1}{h}y'_nv'_1H_1\right)^{-1}$$
(3.1.8)

$$y'_{n+1} = \left(y'_n + \frac{1}{h}w'_1k_1\right)\left(1 - \frac{1}{h}y'_nv'_1H_1\right)$$

Expanding the binomial and re-arranging also gives

$$y'_{n+1} = y'_n + \frac{1}{h}w'_1K_1 - \left(\frac{1}{h}y'_n^2v'_1 + \frac{1}{h^2}y'_nw'_1v'_1K_1\right)H_1$$
(3.1.9)

Now, the Taylor's series expansion of y_{n+1} about x_n is given as

$$y_{n+1} = y_n + hy'_n + \frac{h^2 y_n''}{2!} + \frac{h^3 y_n'''}{3!} + \frac{h^4 y_n^{iv}}{4!} + \dots$$
 (3.1.10)

and

$$y'_{n+1} = y'_n + hy''_n + \frac{h^2 y_n'''}{2!} + \frac{h^3 y_n^{iv}}{3!} + \dots$$
 (3.1.11)

where

$$y_{n}^{"} = f(x_{n}, y_{n}, y_{n}^{"}) = f_{n}$$

$$y_{n}^{"'} = f_{x} + y' f_{y} + f_{n} f_{y} = \Delta f_{n}$$

$$y_{n}^{"} = f_{xx} + y'_{n}^{2} f_{yy} + f^{2} f_{y'y'} + 2y' f_{n} f_{yy'} + 2f_{n} f_{xy'} + f_{y'} \Delta f_{n}$$

$$y_{n}^{"} = \Delta^{2} f_{n} + f_{y'} \Delta f_{n} + f_{n} f_{y}$$

$$(3.1.12)$$

Since
$$\Delta = \frac{\partial}{\partial x} + y' \frac{\partial}{\partial y} + f_n \frac{\partial}{\partial y'}$$
 (3.1.13)

The Taylor's series of the function of three variables is defined as

$$f(x+h,y+A,y'+B) = \sum_{j=0}^{\infty} \left\{ \frac{1}{j!} \left(h \frac{\partial}{\partial x} + A \frac{\partial}{\partial y} + \frac{\partial}{\partial y'} \right)^{j} f(x',y',y'')' \right\}_{x=x',y=y',y'=y''} (3.1.14)$$

This implies that from equation (3.1.3a), we have

$$K_1 = \frac{h^2}{2} f\left(x_n + c_1 h, y_n + h c_1 y'_n + a_{11} K_1, y'_n + \frac{1}{h} b_{11} K_1\right)$$
(3.1.15)

Expanding (3.1.15) in Taylor's series gives

$$\frac{2}{h^{2}}K_{1} = f_{n} + \left(c_{1}hf_{x} + (hc_{1}y'_{n} + a_{11}K_{1})f_{n} + \frac{1}{h}b_{11}f_{y'}\right)
+ \frac{1}{2!}\left((c_{1}h)^{2}f_{xx} + 2c_{1}h(hc_{1}y'_{n} + a_{11}K_{1})f_{xy} + 2c_{1}h\left(\frac{1}{h}b_{11}K_{1}\right)f_{xy'}
+ (hc_{1}y'_{n} + a_{11}K_{1})^{2}f_{yy} + 2(hc_{1}y'_{n} + a_{11}K_{1})\left(\frac{1}{h}b_{11}K_{1}\right)f_{yy'}
+ \left(\frac{1}{h}b_{11}K_{1}\right)^{2}f_{y'y'}\right) + \dots$$
(3.1.16)

this implies that

$$\begin{split} K_1 &= \frac{h^2}{2} f_n + \frac{h^2}{2} \left(c_1 h f_x + (h c_1 y'_n + a_{11} K_1) f_n + \frac{1}{h} b_{11} f_{y'} \right) \\ &+ \frac{h^2}{4} \left((c_1 h)^2 f_{xx} + 2 c_1 h (h c_1 y'_n + a_{11} K_1) f_{xy} + 2 c_1 h \left(\frac{1}{h} b_{11} K_1 \right) f_{xy'} \right. \\ &+ \left. (h c_1 y'_n + a_{11} K_1)^2 f_{yy} + 2 (h c_1 y'_n + a_{11} K_1) \left(\frac{1}{h} b_{11} K_1 \right) f_{yy'} \right. \\ &+ \left. \left(\frac{1}{h} b_{11} K_1 \right)^2 f_{y'y'} \right) + 0 (h^3) \end{split}$$

Simplifying further and arranging the equation in powers of h gives,

$$K_{1} = \frac{h}{2} \left[b_{11} K_{1} f_{y'} + a_{11} b_{11} K_{1}^{2} f_{yy'} \right]$$

$$+ \frac{h^{2}}{2} \left[f_{n} + a_{11} K_{1} f_{y} + c_{1} b_{11} K_{1} f_{xy'} + a_{11}^{2} K_{1}^{2} f_{yy} + c_{1} y'_{n} b_{11} K_{1} f_{yy'} \right]$$

$$+ \frac{h^{3}}{2} \left[c_{1} f_{x} + c_{1} y'_{n} f_{y} + c_{1} a_{11} K_{1} f_{xy} + c_{1} a_{11} y'_{n} K_{1} f_{yy} \right]$$

$$+ \frac{h^{4}}{4} \left[c_{1}^{2} f_{xx} + c_{1}^{2} y'_{n} f_{xy} + c_{1}^{2} y'_{n}^{2} f_{yy} \right] + 0(h^{5})$$

$$(3.1.17)$$

Equation (3.1.17) is implicit; one cannot proceed by successive substitution. Following *Lambert* (1973), we can assume that the solution for K_1 may be express in the form

$$K_1 = hA_1 + h^2B_1 + h^3C_1 + h^4D_1 + 0(h^5)$$
(3.1.18)

Substituting equation (3.1.18) into (3.1.17) gives

$$K_{1} = \frac{h}{2} \left[b_{11} (hA_{1} + h^{2}B_{1} + h^{3}C_{1}) f_{y'} + a_{11}b_{11} (hA_{1} + h^{2}B_{1})^{2} f_{yy'} \right]$$

$$+ \frac{h^{2}}{2} \left[f_{n} + a_{11} (hA_{1} + h^{2}B_{1}) f_{y} + c_{1}b_{11} (hA_{1} + h^{2}B_{1}) f_{xy'} + a_{11}^{2} (hA_{1})^{2} f_{yy} \right]$$

$$+ c_{1}y'_{n}b_{11} (hA_{1} + h^{2}B_{1}) f_{yy'}$$

$$+ \frac{h^{3}}{2} \left[c_{1}f_{x} + c_{1}y'_{n}f_{y} + c_{1}a_{11} (hA_{1}) f_{xy} + c_{1}a_{11}y'_{n} (hA_{1}) f_{yy} \right]$$

$$+ \frac{h^{4}}{4} \left[c_{1}^{2}f_{xx} + c_{1}^{2}y'_{n}f_{xy} + c_{1}^{2}y'_{n}^{2} f_{yy} \right] + 0(h^{5})$$

$$(3.1.19)$$

On equating powers of h from equation (3.1.17) and (3.1.18), gives

$$A_{1} = 0, B_{1} = \frac{1}{2}f_{n}, C_{1} = \frac{1}{2}\left(c_{1}f_{x} + c_{1}y'_{n}f_{y} + \frac{1}{2}b_{11}f_{n}f_{y}\right) = \frac{1}{2}c_{1}\Delta f_{n}, \text{ since } c_{1} = \frac{1}{2}b_{11}$$

$$D_{1} = \frac{1}{4}\left(c_{1}^{2}\Delta^{2}f_{n} + b_{11}\Delta f_{n}f_{y} + a_{11}f_{n}f_{y}\right)$$

$$(3.1.20)$$

Substituting A_1 , B_1 , C_1 and D_1 into (3.1.18) gives

$$K_{1} = \frac{h^{2}}{2} f_{n} + \frac{h^{3}}{2} c_{1} \Delta f_{n} + \frac{h^{4}}{2} \left(c_{1}^{2} \Delta^{2} f_{n} + b_{11} \Delta f_{n} f_{y'} + a_{11} f_{n} f_{y} \right)$$
(3.1.21)

Similarly, expanding H_1 in Taylor's series about (x_n, z_n, z'_n) , from (3.1.3b), we have

$$\begin{split} \frac{2}{h^2}H_1 &= g_n + \left(d_1hg_x + (hd_1z'_n + \alpha_{11}H_1)g_z + \frac{1}{h}\beta_{11}g_{z'}\right) \\ &+ \frac{1}{2!}\bigg((d_1h)^2g_{xx} + 2d_1h(hd_1y'_n + \alpha_{11}H_1)g_{xz} + 2c_1h\left(\frac{1}{h}\beta_{11}H_1\right)g_{xz'} \\ &+ (hdy'_n + \alpha_{11}H_1)^2g_{zz} + 2(hd_1y'_n + \alpha_{11}H_1)\left(\frac{1}{h}\beta_{11}H_1\right)g_{zz'} \\ &+ \left(\frac{1}{h}\beta_{11}H_1\right)^2g_{z'z'}\bigg) + \ \dots \end{split}$$

Expanding further as in the case of K_1 above gives

$$H_{1} = \frac{h}{2} [\beta_{11} H_{1} g_{z'} + \alpha_{11} \beta_{11} H_{1}^{2} g_{zz'}]$$

$$+ \frac{h^{2}}{2} [g_{n} + \alpha_{11} H_{1} g_{z} + d_{1} \beta_{11} H_{1} g_{xz'} + \alpha_{11}^{2} H_{1}^{2} g_{zz} + d_{1} z'_{n} \beta_{11} H_{1} g_{zz'}]$$

$$+ \frac{h^{3}}{2} [d_{1} g_{x} + d_{1} z'_{n} g_{z} + d_{1} \alpha_{11} H_{1} g_{xz} + d_{1} \alpha_{11} z'_{n} H_{1} g_{zz}]$$

$$+ \frac{h^{4}}{4} [c_{1}^{2} g_{xx} + d_{1}^{2} z'_{n} g_{xz} + d_{1}^{2} z'_{n}^{2} g_{zz}] + 0(h^{5})$$
(3.1.22)

Equation (3.1.22) is also implicit which cannot be proceed by successive substitution. Assuming a solution of the equation is of the form

$$H_1 = hL_1 + h^2 M_1 + h^3 N_1 + h^4 R_1 + 0(h^5)$$
(3.1.23)

Substituting the values of H_1 in (3.1.23) into equation (3.1.22) and equating powers of h of the equation, we can get the following after substitutions:

$$L_{1} = 0, M_{1} = \frac{1}{2}g_{n}, N_{1} = \frac{1}{2}d_{1}\Delta g_{n} \text{ and } R_{1} = \frac{1}{4}(d_{1}^{2}\Delta^{2}g_{n} + \beta_{11}\Delta g_{n}g_{z} + \alpha_{11}g_{n}g_{z})$$

$$(3.1.24)$$

Substituting equation (3.1.24) into equation (3.1.23) gives

$$H_1 = \frac{h^2}{2}g_n + \frac{h^3}{2}d_1\Delta g_n + \frac{h^4}{2}(d_1^2\Delta^2 g_n + \beta_{11}\Delta g_n g_{z'} + \alpha_{11}g_n g_z)$$
(3.1.25)

Using equations (3.1.18) and (3.1.23) into equations (3.1.7) and (3.1.9) respectively gives

$$y_{n+1} = y_n + hy'_n - (y_n^2 v_1 + hy_n y'_n v_1)(h^2 M_1 + h^3 N_1 + h^4 R_1)$$
$$+ [w_1 - y_n v_1 w_1 (h^2 M_1 + h^3 N_1 + h^4 R_1)](h^2 B_1 + h^3 C_1 + h^4 D_1)$$

Expanding the brackets and re-arranging in powers of h gives

$$y_{n+1} = y_n + hy_n' + h^2(w_1B_1 - y_n^2v_1M_1) + h^3(w_1C_1 - y_n^2v_1N_1 - y_ny_n'v_1M_1) + 0(h^4)$$
(3.1.26)

Also for y'_{n+1} gives

$$y'_{n+1} = y'_n + \frac{1}{h}w'_1(h^2B_1 + h^3C_1 + h^4D_1)$$
$$-\left[\frac{1}{h}y'^2_nv'_n + \frac{1}{h^2}y'_nw'_1v'_n(h^2B_1 + h^3C_1 + h^4D_1)\right](h^2M_1 + h^3N_1 + h^4R_1)$$

Expanding the brackets and re-arrange in powers of h gives

$$y'_{n+1} = y'_n + h(w'_1B_1 - y'^2_nv'_nM_1) + h^2(w'_1C_1 - y'^2_nv'_nN_1 - y'_nw'_1v'_nB_1M_1) + h^3(w'_1D_1 - y'^2_nv'_nR_1 - y'_nw'_1v'_nB_1N_1 - y'_nw'_1v'_nC_1M_1) + 0(h^4) \dots (3.1.27)$$

Comparing the corresponding powers in h of equations (3.1.26) and (3.1.27) with equations (3.1.10) and (3.1.11) we obtain

$$\frac{1}{2}w_{1}f_{n} - \frac{1}{2}y'_{n}v_{1}g_{n} = \frac{1}{2}f_{n}$$

$$\frac{1}{2}w_{1}wc_{1}\Delta f_{n} - \frac{1}{2}y'_{n}^{2}v_{1}d_{1}\Delta g_{n} - \frac{1}{2}y_{n}y'_{n}v_{1}g_{n} = \frac{1}{6}\Delta f_{n}$$

$$\frac{1}{2}w'_{1}f_{n} - \frac{1}{2}y'_{n}^{2}v'_{1}g_{n} = f_{n}$$

$$\frac{1}{2}w'_{1}c_{1}\Delta f_{n} - \frac{1}{2}y'_{n}^{2}v'_{1}d_{1}\Delta g_{n} - \frac{1}{2}y'_{n}w'_{1}v'_{1}f_{n}(\frac{1}{2}g_{n}) = \frac{1}{2}\Delta f_{n}$$
(3.1.28)

(By using the equations in (3.1.20) and (3.1.24))

Since from equation (3.7)

$$g_{n} = -\frac{f_{n}}{y_{n}^{2}}, \quad g_{x} = -\frac{f_{x}}{y_{n}^{2}}, \quad g_{z} = -2\frac{f_{n}}{y_{n}} + f_{y}, \quad g_{z'} = -2\frac{f_{n}}{y_{n}} + f_{y'}, \quad z'_{n} = -\frac{y'_{n}}{y_{n}^{2}}$$
and
$$\Delta g_{n} = g_{n} + z'_{n}g_{z} + g_{n}g_{z'}$$

$$(3.1.29)$$

Using those equations into equation (3.1.28), we get the following simultaneous equations

$$w_{1} + v_{1} = 1$$

$$w_{1}c_{1} + v_{1}d_{1} = \frac{1}{3}$$

$$w'_{1} + v'_{1} = 2$$

$$w'_{1}c_{1} + v'_{1}d_{1} = 1$$

$$(3.1.30)$$

Equation (3.1.30) has (4) equations with (6) unknowns; there will not be a unique solution for (3.1.30). There will be a family of one-stage scheme of order four.

i. Choosing the parameters

$$w_1 = \frac{1}{3}$$
, $v_1 = \frac{2}{3}$, $c_1 = a_{11} = b_{11} = 0$, $w_1' = 0$, $v_1' = 2$, $d_1 = \alpha_{11} = \frac{1}{2}$, $\beta_{11} = 1$ arbitrarily the following scheme is obtain.

$$y_{n+1} = \frac{hy_n' + \frac{1}{3}K_1}{1 + \frac{2}{3}y_nH_1}$$
 (3.1.32)

and

$$y'_{n+1} = \frac{y'_n}{1 + \frac{2}{h}H_1} \tag{3.1.33}$$

where

$$K_{1} = \frac{h^{2}}{2} f(x_{n}, y_{n}, y'_{n})$$

$$H_{1} = \frac{h^{2}}{2} f\left(x_{n} + \frac{1}{2}h, z_{n} + \frac{1}{2}hz'_{n} + \frac{1}{2}H_{1}, z'_{n} + \frac{1}{h}H_{1}\right) , \text{ since } d_{1} = \alpha_{11} = \frac{1}{2}\beta_{11}$$

ii. Choosing the parameters

From (3.1.30) setting

$$w_1 = v_1 = \frac{1}{2}, \quad c_1 = a_{11} = \frac{1}{2}, \quad d_1 = \alpha_{11} = \frac{1}{6}, \quad w_1' = 2, \quad v_1' = 0, \quad b_{11} = 1, \quad \beta_{11} = \frac{1}{3}$$

Then,

$$y_{n+1} = \frac{y_n + hy_n' + \frac{1}{2}K_1}{1 + \frac{1}{2}y_n H_1}$$
 (3.1.34)

and

$$y'_{n+1} = y'_n + \frac{2}{h}K_1 \tag{3.1.35}$$

where

$$K_1 = \frac{h^2}{2} f\left(x_n + \frac{1}{2}h, y_n + \frac{1}{2}hy_n' + \frac{1}{2}K_1, y_n' + \frac{1}{h}K_1\right)$$
, since $c_1 = a_{11} = \frac{1}{2}b_{11}$

and

$$H_1 = \frac{h^2}{2} g \left(x_n + \frac{1}{6} h, z_n + \frac{1}{6} h z'_n + \frac{1}{6} H_1, z'_n + \frac{1}{3h} H_1 \right), \text{ since } d_1 = \alpha_{11} = \frac{1}{2} \beta_{11}$$

3.2 DERIVATION OF TWO - STAGE SCHEME

From equations (3.1), (3.2), (3.3) and (3.4), setting s = 2, we can have

$$y_{n+1} = \frac{y_n + hy'_n + w_1K_1 + w_2K_2}{1 + y_n(v_1H_1 + v_2H_2)}$$
(3.2.1)

and

$$y'_{n+1} = \frac{y_n + \frac{1}{h}(w'_1 K_1 + w'_2 K_2)}{1 + \frac{1}{h}y'_n(v'_1 H_1 + v'_2 H_2)}$$
(3.2.2)

where

$$k_2 = \frac{h^2}{2} f\left(x_n + c_2 h, y_n + h c_2 y'_n + a_{21} K_1 + a_{22} K_2, y'_n + \frac{1}{h} (b_{21} K_1 + b_{22} K_2)\right)$$
(3.2.3a)

and

$$H_{2} = \frac{h^{2}}{2}g\left(x_{n} + d_{2}h, z_{n} + hd_{2}z'_{n} + \alpha_{21}H_{1} + \alpha_{22}H_{2}, z'_{n} + \frac{1}{h}(\beta_{21}H_{1} + \beta_{22}H_{2})\right)$$
(3.2.3b)

with constraints

$$c_2 = a_{21} + a_{22} = \frac{1}{2}(b_{21} + b_{22})$$
 and $d_2 = a_{21} + a_{22} = \frac{1}{2}(\beta_{21} + \beta_{22})$ (3.2.4)

Where c_2 , a_{21} , a_{22} , b_{21} , b_{22} d_2 , a_{21} , a_{22} , a_{21} , a_{22} , a_{21} , a_{22} , a_{22} , a_{21} , a_{22} , a_{22

Now by adopting a binomial expansion on equations (3.2.1) and (3.2.2), gives This implies that from equation (3.2.1), we have

$$y_{n+1} = (y_n + hy'_n + w_1K_1 + w_2K_2)(1 + y_n(v_1H_1 + v_2H_2))^{-1}$$

$$y_{n+1} = (y_n + hy'_n + w_1K_1 + w_2K_2)(1 - y_n(v_1H_1 + v_2H_2)) - \text{Higher order terms}$$

Expanding, re-arranging and ignoring the higher order terms, gives

$$y_{n+1} = y_n + hy'_n + w_1K_1 + w_2K_2 - y_n^2(v_1H_1 + v_2H_2) - hy_ny'_n(v_1H_1 + v_2H_2) - y_n(w_1K_1 + w_2K_2)(v_1H_1 + v_2H_2)$$

$$(3.2.5)$$

also

$$y'_{n+1} = \left(y'_n + \frac{1}{h}(w'_1K_1 + w'_2K_2)\right)\left(1 - \frac{1}{h}y'_n(v'_1H_1 + v'_2H_2)\right)^{-1}$$

Expanding and re-arranging also gives

$$y'_{n+1} = y'_n - \frac{1}{h} y'_n^2 (v'_1 H_1 + v'_2 H_2) + \frac{1}{h} (w'_1 K_1 + w'_2 K_2)$$
$$- \frac{1}{h^2} y'_n (w'_1 K_1 + w'_2 K_2) (v'_1 H_1 + v'_2 H_2)$$
(3.2.6)

From equations (3.2.3a) and (3.2.3b), the Taylor series expansion of K_2 about (x_n, y_n, y'_n) is given as

$$\frac{2}{h^{2}}K_{2} = f_{n} + \left(c_{2}hf_{x} + (hc_{2}y'_{n} + a_{21}K_{1} + a_{22}K_{2})f_{n} + \frac{1}{h}(b_{21}K_{1} + b_{22}K_{2})f_{yy}\right)
+ \frac{1}{2!}\left((c_{2}h)^{2}f_{xx} + 2c_{2}h(hc_{1}y'_{n} + a_{21}K_{1} + a_{22}K_{2})f_{xy} + 2c_{2}(b_{21}K_{1} + b_{22}K_{2})f_{xy} + (hc_{2}y'_{n} + a_{21}K_{1} + a_{22}K_{2})^{2}f_{yy} + \frac{2}{h}(a_{21}K_{1} + a_{22}K_{2})(b_{21}K_{1} + b_{22}K_{2})f_{yy} + \left(\frac{1}{h}(b_{21}K_{1} + b_{22}K_{2})\right)^{2}f_{yyy} + \dots$$

$$+ \dots$$

Expanding the brackets and re-arranging in powers of h gives

$$\begin{split} \frac{2}{h^2}K_2 &= f_n + (a_{21}K_1 + a_{22}K_2)f_y \\ &\quad + h\big[c_2f_x + c_2y_n'f_y + c_2(a_{21}K_1 + a_{22}K_2)f_{xy} + 2c_2y_n' \ (a_{21}K_1 + a_{22}K_2)f_{yy}\big] \\ &\quad + \frac{h^2}{2}\big[c_2^2f_{xx} + 2c_2y_n'f_{xy} + c_2^2y_n'^2f_{yy}\big] \\ &\quad + \frac{1}{2}\big[2c_2(b_{21}K_1 + b_{22}K_2)f_{xy} + (a_{21}K_1 + a_{22}K_2)^2f_{yy} \\ &\quad + 2c_2y_n'(b_{21}K_1 + b_{22}K_2)f_{yy}\big] + \frac{2}{h}(a_{21}K_1 + a_{22}K_2)(b_{21}K_1 + b_{22}K_2)f_{yy} \\ &\quad + \frac{1}{h}(b_{21}K_1 + b_{22}K_2)f_{yy} + \frac{1}{h^2}(b_{21}K_1 + b_{22}K_2)^2f_{yyy} + 0(h^3) \end{split}$$

This implies that

$$K_{2} = \frac{h^{2}}{2} \left[f_{n} + (a_{21}K_{1} + a_{22}K_{2})f_{y} + c_{2}(b_{21}K_{1} + b_{22}K_{2})f_{xy'} + (a_{21}K_{1} + a_{22}K_{2})^{2} f_{yy} \right.$$

$$+ c_{2}y'_{n}(b_{21}K_{1} + b_{22}K_{2})f_{yy'} \left]$$

$$+ \frac{h^{3}}{2} \left[c_{2}f_{x} + c_{2}y'_{n}f_{y} + c_{2}(a_{21}K_{1} + a_{22}K_{2})f_{xy} + c_{2}y'_{n}(a_{21}K_{1} + a_{22}K_{2})f_{yy} \right]$$

$$+ \frac{h^{4}}{4} \left[c_{2}^{2}f_{xx} + 2c_{2}y'_{n}f_{xy'} + c_{2}^{2}y'_{n}^{2}f_{yy} \right]$$

$$+ h(a_{21}K_{1} + a_{22}K_{2})(b_{21}K_{1} + b_{22}K_{2})f_{yy'} + \frac{h}{2}(b_{21}K_{1} + b_{22}K_{2})f_{y'}$$

$$+ \frac{1}{2}(b_{21}K_{1} + b_{22}K_{2})^{2}f_{y'y'} + 0(h^{5})$$
(3.2.7)

Equation (3.2.7) is implicit, which cannot be proceed by successive substitutions. We assume a solution for K_2 which may be expressed as

$$K_2 = h^2 B_2 + h^3 C_2 + h^4 D_2 + 0(h^5)$$
(3.2.8)

Substituting equation (3.2.8) into (3.2.7) gives

$$\begin{split} K_2 &= \frac{h^2}{2} \Big[f_n + \Big(a_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + a_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) f_y \\ &+ c_2 \Big(b_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + b_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) f_{xy}, \\ &+ \Big(a_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + a_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big)^2 f_{yy} \\ &+ c_2 y_n' \Big(b_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + b_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) f_{yy}, \Big] \\ &+ \frac{h^3}{2} \Big[c_2 f_x + c_2 y_n' f_y \\ &+ c_2 \Big(a_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + a_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) f_{xy} \\ &+ c_2 y_n' \Big(a_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + a_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) f_{yy} \Big] \\ &+ \frac{h^4}{4} \Big[c_2^2 f_{xx} + 2 c_2 y_n' f_{xy}, + c_2^2 y_n'^2 f_{yy} \Big] \\ &+ h \Big(a_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) \Big) \\ &+ a_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) \Big(b_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) \\ &+ b_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) f_{yy}, \\ &+ \frac{h}{2} \Big(b_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + b_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big) f_{yy}, \\ &+ \frac{1}{2} \Big(b_{21} (h^2 B_1 + h^3 C_1 + h^4 D_1) + b_{22} (h^2 B_2 + h^3 C_2 + h^4 D_2) \Big)^2 f_{yyy}, + \cdots \end{aligned}$$

Expanding and re-arranging gives

$$K_{2} = \frac{h^{2}}{2} f_{n} + \frac{h^{3}}{2} \left[c_{2} f_{x} + c_{2} y_{n}' f_{y} + (b_{21} B_{1} + b_{22} B_{2}) f_{y'} \right]$$

$$+ \frac{h^{4}}{4} \left[c_{2}^{2} f_{xx} + 2 c_{2} y_{n}' f_{xy'} + c_{2}^{2} y_{n}'^{2} f_{yy} + 2 (b_{21} B_{1} + b_{22} B_{2})^{2} f_{yy'} \right]$$

$$+ 2 c_{2} (b_{21} B_{1} + b_{22} B_{2}) f_{xy'} + c_{2} y_{n}' (b_{21} B_{1} + b_{22} B_{2}) f_{yy'}$$

$$+ 2 (a_{21} B_{1} + a_{22} B_{2}) f_{y} + 2 (b_{21} C_{1} + b_{22} C_{2}) f_{yy'} \right] + 0 (h^{5})$$

$$(3.2.9)$$

On equating powers of h from equation (3.2.9) and (3.2.8), gives

$$B_{2} = \frac{1}{2} f_{n}$$

$$C_{2} = \frac{1}{2} \left(c_{2} f_{x} + c_{2} y_{n}' f_{y} + (b_{21} B_{1} + b_{22} B_{2}) f_{y'} \right) = \frac{1}{2} c_{2} \Delta f_{n}$$

$$D_{2} = \frac{1}{4} \left(c_{2}^{2} \Delta^{2} f_{n} + 2 c_{2} \Delta f_{n} f_{y'} + c_{2} f_{n} f_{y} \right)$$

$$(3.2.10)$$

since
$$\left(c_2 = a_{21} + a_{22} = \frac{1}{2}(b_{21} + b_{22})\right)$$

then

$$K_2 = \frac{h^2}{2} f_n + \frac{h^3}{2} c_2 \Delta f_n + \frac{h^4}{4} \left(c_2^2 \Delta^2 f_n + 2c_2 \Delta f_n f_{y'} + c_2 f_n f_y \right) + 0(h^5)$$
 (3.2.11)

In a similar manner

$$H_2 = h^2 M_2 + h^3 N_2 + h^4 R_2 + 0(h^5) (3.2.12)$$

Where

$$M_{2} = \frac{1}{2}g_{n}$$

$$N_{2} = \frac{1}{2}d_{2}\Delta g f_{n}$$

$$R_{2} = \frac{1}{4}(d_{2}^{2}\Delta^{2}g_{n} + 2d_{2}\Delta g_{n}g_{z'} + d_{2}g_{n}g_{z})$$

$$(3.2.13)$$

And also,

$$H_2 = \frac{h^2}{2}g_n + \frac{h^3}{2}d_2\Delta g f_n + \frac{h^4}{4}(d_2^2\Delta^2 g_n + 2d_2\Delta g_n g_{z'} + d_2g_n g_z) + 0(h^5)$$
 (3.2.14)

Substituting equations (3.1.20), (3.1.23), (3.2.8) and (3.2.12) into equations (3.2.5) and (3.2.6), re – arranging and compare the resulting equation with equations (3.1.7) and (3.1.8) gives the following

$$w_{1}B_{1} + w_{2}B_{2} - y_{n}^{2}(v_{1}M_{1} + v_{2}M_{2}) = \frac{1}{2}f_{n}$$

$$w_{1}C_{1} + w_{2}C_{2} - y_{n}^{2}(v_{1}M_{1} + v_{2}M_{2}) = \frac{1}{6}\Delta f_{n}$$

$$w'_{1}B_{1} + w'_{2}B_{2} - y_{n}^{2}(v'_{1}M_{1} + v'_{2}M_{2}) = f_{n}$$

$$w'_{1}C_{1} + w'_{2}C_{2} - y_{n}^{\prime 2}(v'_{1}M_{1} + v'_{2}M_{2}) = \frac{1}{2}\Delta f_{n}$$

$$(3.2.15)$$

Substituting the values of B_1 , B_2 , C_1 , C_2 , M_1 , M_2 , N_2 , and N_2 , we have the following non – liner differential equations

$$w_{1} + w_{2} + v_{1} + v_{2} = 1$$

$$w_{1}c_{1} + w_{2}c_{2} + v_{1}d_{1} + v_{2}d_{2} = \frac{1}{3}$$

$$w'_{1} + w'_{2} + v'_{1} + v'_{2} = 2$$

$$w'_{1}c_{1} + w'_{1}c_{1} + v'_{1}d_{1} + v'_{2}d_{2} = 1$$

$$(3.2.16)$$

Which are four equations with twelve unknowns from (3.2.16). This implies that there are families of solutions.

1. From (3.2.16) setting
$$w_1 = w_2 = 0$$
, $v_1 = \frac{1}{4}$, $v_2 = \frac{3}{4}$, $w_1' = w_2' = 0$, $v_1' = v_2' = 1$, $d_1 = \frac{1}{6}$, $d_2 = \frac{5}{6}$

Then we have,

$$y_{n+1} = \frac{y_n + hy_n'}{1 + y_n^{\frac{1}{2}}(H_1 + 3H_2)}$$
(3.2.17)

and

$$y'_{n+1} = \frac{y'_n}{1 + \frac{1}{h} y'_n (H_1 + H_2)}$$
 (3.2.18)

where

$$H_{1} = \frac{h^{2}}{2} g \left(x_{n} + \frac{1}{6} h, z_{n} + \frac{1}{6} h z'_{n} + \frac{1}{6} H_{1}, z'_{n} + \frac{1}{3h} H_{1} \right)$$

$$H_{2} = \frac{h^{2}}{2} g \left(x_{n} + \frac{5}{6} h, z_{n} + \frac{5}{6} h z'_{n} + \frac{1}{2} H_{1} + \frac{1}{3} H_{2}, z'_{n} + \frac{1}{h} \left(H_{1} + \frac{2}{3} H_{2} \right) \right)$$

2. Also from (3.2.16) setting $w_1 = w_2 = \frac{1}{8}$, $v_1 = v_2 = \frac{3}{8}$, $c_1 = c_2 = d_1 = d_2 = \frac{1}{3}$ and $w_1' = w_2' = v_1' = v_2' = \frac{1}{2}$

Then we have,

$$y_{n+1} = \frac{y_n + hy_n' + \frac{1}{8}(K_1 + K_2)}{1 + y_n \frac{3}{8}(H_1 + H_2)}$$
(3.2.19)

and

$$y'_{n+1} = \frac{y'_n + \frac{1}{2h}(K_1 + K_2)}{1 + \frac{1}{2h}y'_n(H_1 + H_2)}$$
(3.2.20)

where

$$K_{1} = \frac{h^{2}}{2} f\left(x_{n} + \frac{1}{3}h, y_{n} + \frac{1}{3}hy'_{n} + \frac{1}{3}K_{1}, y'_{n} + \frac{2}{3h}K_{1}\right)$$

$$K_{2} = \frac{h^{2}}{2} f\left(x_{n} + \frac{1}{3}h, y_{n} + \frac{1}{3}hy'_{n} + \frac{1}{3}(K_{1} + K_{2}), y'_{n} + \frac{2}{3h}(K_{1} + K_{2})\right)$$

and

$$H_{1} = \frac{h^{2}}{2} g \left(x_{n} + \frac{1}{3} h, z_{n} + \frac{1}{3} h z'_{n} + \frac{1}{3} H_{1}, z'_{n} + \frac{2}{3h} H_{1} \right)$$

$$H_{2} = \frac{h^{2}}{2} g \left(x_{n} + \frac{1}{3} h, z_{n} + \frac{1}{3} h z'_{n} + \frac{1}{3} (H_{1} + H_{2}), z'_{n} + \frac{2}{3h} (H_{1} + H_{2}) \right)$$

3.3 DERIVATION OF THREE - STAGE SCHEME

From equations (3.1), (3.2), (3.3) and (3.4), setting s = 3, we can have

$$y_{n+1} = \frac{y_n + hy_n' + w_1K_1 + w_2K_2 + w_3K_3}{1 + y_n(v_1H_1 + v_2H_2 + v_3H_3)}$$
(3.3.1)

and

$$y'_{n+1} = \frac{y_n + \frac{1}{h}(w'_1 K_1 + w'_2 K_2 + w'_3 K_3)}{1 + \frac{1}{h}y'_n(v'_1 H_1 + v'_2 H_2 + v'_3 H_3)}$$
(3.3.2)

where

$$K_{3} = \frac{h^{2}}{2} f\left(x_{n} + c_{3}h, y_{n} + hc_{3}y'_{n} + a_{31}K_{1} + a_{32}K_{2} + a_{33}K_{3}, y'_{n} + \frac{1}{h}(b_{31}K_{1} + b_{32}K_{2} + b_{33}K_{3})\right),$$

$$(3.3.3a)$$

and

$$H_{3} = \frac{h^{2}}{2} g \left(x_{n} + d_{3}h, z_{n} + hd_{3}z'_{n} + \alpha_{31}H_{1} + \alpha_{32}H_{2} + \alpha_{33}H_{3}, z'_{n} + \frac{1}{h} (\beta_{31}H_{1} + \beta_{32}H_{2} + \beta_{33}H_{3}) \right)$$

$$(3.2.3b)$$

with constraints

$$c_{3} = a_{31} + a_{32} + a_{33} = \frac{1}{2}(b_{31} + b_{32} + b_{33})$$

$$d_{3} = \alpha_{31} + \alpha_{32} + \alpha_{33} = \frac{1}{2}(\beta_{31} + \beta_{32} + \beta_{33})$$

$$(3.3.4)$$

where c_3 , a_{31} , a_{32} , a_{33} b_{31} , b_{32} , b_{33} , d_3 , a_{31} , a_{23} , a_{23

Now by adopting a binomial expansion on equations (3.3.1) and (3.3.2), gives This implies that from equation (3.3.1), we have

$$y_{n+1} = (y_n + hy'_n + w_1K_1 + w_2K_2 + w_3K_3)(1 + y_n(v_1H_1 + v_2H_2 + v_3H_3))^{-1}$$

$$y_{n+1} = (y_n + hy'_n + w_1K_1 + w_2K_2 + w_3K_3)(1 - y_n(v_1H_1 + v_2H_2 + v_3H_3)) - \text{Higher}$$
order terms

Expanding, re-arranging and ignoring the higher order terms, gives

$$y_{n+1} = y_n + hy'_n + w_1K_1 + w_2K_2 + w_3K_3 - y_n^2(v_1H_1 + v_2H_2 + v_3H_3)$$

$$- hy_ny'_n(v_1H_1 + v_2H_2 + v_3H_3)$$

$$- y_n[(w_1K_1 + w_2K_2 + w_3K_3)(v_1H_1 + v_2H_2 + v_3H_3)]$$

$$+ v_3H_3)]$$
(3.3.5)

also

$$y'_{n+1} = \left(y'_n + \frac{1}{h}(w'_1K_1 + w'_2K_2 + w'_3K_3)\right) \left(1 - \frac{1}{h}y'_n(v'_1H_1 + v'_2H_2 + v'_3K_3)\right)^{-1}$$

Expanding and re-arranging also gives

$$y'_{n+1} = y'_n - \frac{1}{h} y'_n^2 (v'_1 H_1 + v'_2 H_2 + v'_3 K_3) + \frac{1}{h} (w'_1 K_1 + w'_2 K_2 + w'_3 K_3)$$
$$- \frac{1}{h^2} y'_n (w'_1 K_1 + w'_2 K_2 + w'_3 K_3) (v'_1 H_1 + v'_2 H_2$$
$$+ v'_3 K_3)$$
(3.3.6)

Now, expanding K_3 Taylor series gives

$$\begin{split} \frac{2}{h^2}K_3 &= f_n + \left(c_3hf_x + (hc_3y'_n + a_{31}K_1 + a_{32}K_2 + a_{33}K_3)f_y \right. \\ &\quad + \frac{1}{h}(b_{31}K_1 + b_{32}K_2 + b_{33}K_3)f_{y'} \Big) \\ &\quad + \frac{1}{2!}\left((c_3h)^2f_{xx} + 2c_3h(hc_1y'_n + a_{31}K_1 + a_{32}K_2 + a_{33}K_3)f_{xy} \right. \\ &\quad + 2c_2(b_{31}K_1 + b_{32}K_2 + b_{33}K_3)f_{xy'} \\ &\quad + (hc_2y'_n + a_{31}K_1 + a_{32}K_2 + a_{33}K_3)^2f_{yy} \\ &\quad + \frac{2}{h}(a_{31}K_1 + a_{32}K_2 + a_{33}K_3)(b_{31}K_1 + b_{32}K_2 + b_{33}K_3)f_{yy'} \\ &\quad + \frac{1}{h^2}\left((b_{31}K_1 + b_{32}K_2 + b_{33}K_3)\right)^2f_{y'y'}\right) + \dots \end{split}$$

Expanding the brackets and re-arranging in powers of h gives

$$\begin{split} \frac{2}{h^2}K_3 &= f_n + (a_{31}K_1 + a_{32}K_2 + a_{33}K_3)f_y \\ &\quad + h \big[c_3f_x + c_3y_n'f_y + c_3(a_{31}K_1 + a_{32}K_2 + a_{33}K_3)f_{xy} \\ &\quad + 2c_3y_n' \left(a_{31}K_1 + a_{32}K_2 + a_{33}K_3 \right)f_{yy} \big] + \frac{h^2}{2} \big[c_3{}^2f_{xx} + 2c_3y_n'f_{xy} + c_3^2y_n'^2f_{yy} \big] \\ &\quad + \frac{1}{2} \big[2c_3(b_{31}K_1 + b_{32}K_2 + b_{33}K_3)f_{xy} + (a_{31}K_1 + a_{32}K_2 + a_{33}K_3)^2f_{yy} \\ &\quad + 2c_3y_n'(b_{31}K_1 + b_{32}K_2 + b_{33}K_3)f_{yy} \big] \\ &\quad + \frac{2}{h} (a_{31}K_1 + a_{32}K_2 + a_{33}K_3)(b_{21}K_1 + b_{22}K_2)f_{yy} \\ &\quad + \frac{1}{h} (b_{31}K_1 + b_{32}K_2 + b_{33}K_3)f_{y} + \frac{1}{h^2} (b_{31}K_1 + b_{32}K_2 + b_{33}K_3)^2f_{yy} \\ &\quad + 0(h^3) \end{split}$$

this implies that

$$K_{3} = \frac{h^{2}}{2} \left[f_{n} + (a_{31}K_{1} + a_{32}K_{2} + a_{33}K_{3}) f_{y} + 2c_{3}(b_{31}K_{1} + b_{32}K_{2} + b_{33}K_{3}) f_{xy'} \right.$$

$$+ (a_{31}K_{1} + a_{32}K_{2} + a_{33}K_{3})^{2} f_{yy} + 2c_{3}y'_{n}(b_{31}K_{1} + b_{32}K_{2} + b_{33}K_{3}) f_{yy'} \right]$$

$$+ \frac{h^{3}}{2} \left[c_{3}f_{x} + c_{3}y'_{n}f_{y} + c_{3}(a_{31}K_{1} + a_{32}K_{2} + a_{33}K_{3}) f_{xy} \right.$$

$$+ c_{3}y'_{n}(a_{31}K_{1} + a_{32}K_{2} + a_{33}K_{3}) f_{yy} \right] + \frac{h^{4}}{4} \left[c_{3}^{2} f_{xx} + 2c_{3}y'_{n} f_{xy} + c_{3}^{2}y'_{n}^{2} f_{yy} \right]$$

$$+ h(a_{31}K_{1} + a_{32}K_{2} + a_{33}K_{3})(b_{31}K_{1} + b_{32}K_{2} + b_{33}K_{3}) f_{yy'}$$

$$+ \frac{h}{2}(b_{31}K_{1} + b_{32}K_{2} + b_{33}K_{3}) f_{y'} + \frac{1}{2}(b_{31}K_{1} + b_{32}K_{2} + b_{33}K_{3})^{2} f_{y'y'}$$

$$+ 0(h^{5})$$

$$(3.3.7)$$

Equation (3.3.7) is implicit, which cannot be proceed by successive substitutions. We assume a solution for K_2 which may be expressed as

$$K_3 = h^2 B_3 + h^3 C_3 + h^4 D_3 + 0(h^5)$$
(3.3.8)

Substituting the values of K_1 , K_2 and K_3 into equation (3.2.7) and re – arranging in powers of h

Expanding and re-arranging gives

$$K_{3} = \frac{h^{2}}{2} f_{n} + \frac{h^{3}}{2} \left[c_{3} f_{x} + c_{3} y_{n}' f_{y} + \frac{1}{2} (b_{31} B_{1} + b_{32} B_{2} + b_{33} B_{3}) f_{y'} \right]$$

$$+ \frac{h^{4}}{4} \left[c_{3}^{2} f_{xx} + 2 c_{3} y_{n}' f_{xy'} + c_{3}^{2} y_{n}'^{2} f_{yy} + 2 (b_{31} B_{1} + b_{32} B_{2} + b_{33} B_{3})^{2} f_{yy'} + 2 (b_{31} B_{1} + b_{32} B_{2} + b_{33} B_{3}) f_{xy'} + c_{2} y_{n}' (b_{31} B_{1} + b_{32} B_{2} + b_{33} B_{3}) f_{yy'} + 2 (b_{31} B_{1} + b_{32} B_{2} + b_{33} B_{3}) f_{yy'} + 2 (b_{31} C_{1} + b_{32} C_{2} + b_{33} C_{3}) f_{y'} \right]$$

$$+ 0 (h^{5})$$

$$(3.3.9)$$

On equating powers of h from equation (3.3.9) and (3.3.8), gives

$$B_{3} = \frac{1}{2} f_{n}$$

$$C_{3} = \frac{1}{2} \left(c_{3} f_{x} + c_{3} y_{n}' f_{y} + (b_{31} B_{1} + b_{32} B_{2} + b_{33} B_{3}) f_{y'} \right) = \frac{1}{2} c_{3} \Delta f_{n}$$

$$D_{2} = \frac{1}{4} \left(c_{3}^{2} \Delta^{2} f_{n} + (b_{31} C_{1} + b_{32} C_{2} + b_{33} C_{3}) \Delta f_{n} f_{y'} + c_{3} f_{n} f_{y} \right)$$

$$(3.3.10)$$

then

$$K_{2} = \frac{h^{2}}{2} f_{n} + \frac{h^{3}}{2} c_{3} \Delta f_{n} + \frac{h^{4}}{4} \left(c_{3}^{2} \Delta^{2} f_{n} + (b_{31} C_{1} + b_{32} C_{2} + b_{33} C_{3}) \Delta f_{n} f_{y'} + c_{2} f_{n} f_{y} \right)$$

$$+ 0(h^{5})$$
(3.3.11)

In a similar manner

$$H_3 = h^2 M_3 + h^3 N_3 + h^4 R_3 + 0(h^5) (3.3.12)$$

Where

$$M_{3} = \frac{1}{2}g_{n}$$

$$N_{3} = \frac{1}{2}d_{3}\Delta g f_{n}$$

$$R_{3} = \frac{1}{4}(d_{3}^{2}\Delta^{2}g_{n} + (\beta_{31}d_{1} + \beta_{32}d_{2} + \beta_{33}d_{3})\Delta g_{n}g_{z'} + d_{3}g_{n}g_{z})$$

$$(3.3.13)$$

And also,

$$H_{3} = \frac{h^{2}}{2}g_{n} + \frac{h^{3}}{2}d_{3}\Delta g f_{n} + \frac{h^{4}}{4}(d_{3}^{2}\Delta^{2}g_{n} + (\beta_{31}d_{1} + \beta_{32}d_{2} + \beta_{33}d_{3})\Delta g_{n}g_{z'} + d_{3}g_{n}g_{z}) + 0(h^{5})$$
(3.3.14)

Substituting equations (3.1.15), (3.1.21), (3.2.8), (3.3.8), (3.2.12) and (3.3.12) into equations (3.2.5), (3.2.6), (3.3.5) and (3.3.6) re – arranging and compare the resulting equation with equations (3.1.7) and (3.1.8) gives the following

$$\begin{split} w_1B_1 + w_2B_2 + w_3B_3 - y_n^2(v_1M_1 + v_2M_2 + v_3M_3) &= \frac{1}{2}f_n \\ w_1C_1 + w_2C_2 + w_3B_3 - y_n^2(v_1N_1 + v_2N_2 + v_3N_3) &= \frac{1}{6}\Delta f_n \\ w_1D_1 + w_2D_2 + w_3D_3 - y_n^2(v_1R_1 + v_2R_2 + v_3R_3) &= \frac{1}{24}\left(\Delta^2 f_n + \Delta f_n f_{y'} + f_n f_y\right) \\ w_1'B_1 + w_2'B_2 + w_3'B_3 - y_n^2(v_1'M_1 + v_2'M_2 + v_3'M_3) &= f_n \\ w_1'C_1 + w_2'C_2 + w_3'C_3 - y_n'^2(v_1'N_1 + v_2'N_2 + v_3'N_3) &= \frac{1}{2}\Delta f_n \end{split}$$

$$w_1'D_1 + w_2'D_2 + w_3'D_3 - y_n'^2(v_1'R_1 + v_2'R_2 + v_3'R_3) = \frac{1}{6}(\Delta^2 f_n + \Delta f_n f_{y'} + f_n f_y)$$

Substituting the values of B_1 , B_2 , B_3 , C_1 , C_2 , C_3 , M_1 , M_2 , M_3 , N_1 , N_2 , N_3 , R_1 , R_2 , and R_3 in the above equation and re-arranging, we have the following sets of non – liner differential equations

$$\begin{split} w_1 + w_2 + w_3 + v_1 + v_2 + v_3 &= 1 \\ w_1c_1 + w_2c_2 + w_3c_3 + v_1d_1 + v_2d_2 + v_3d_3 &= \frac{1}{3} \\ w_1c_1^2 + w_2c_2^2 + w_3c_3^2 + v_1d_1^2 + v_2d_2^2 + v_3d_3^2 &= \frac{1}{6} \\ w_1(c_1b_{11} + c_2b_{12} + c_3b_{13}) + w_2(c_1b_{21} + c_2b_{22} + c_3b_{23}) + w_3(c_1b_{31} + c_2b_{32} + c_3b_{33}) + v_1(\beta_{11}d_1 + \beta_{12}d_2 + \beta_{13}d_3) + v_2(\beta_{21}d_1 + \beta_{22}d_2 + \beta_{23}d_3) + v_3(\beta_{31}d_1 + \beta_{32}d_2 + \beta_{33}d_3) &= \frac{1}{6} \\ w_1' + w_2' + w_3' + v_1' + v_2' + v_3' &= 2 \\ w_1'c_1 + w_2'c_2 + w_3'c_3 + v_1'd_1 + v_2'd_2 + v_3'd_3 &= 1 \\ w_1'c_1^2 + w_2'c_2^2 + w_3'c_3^2 + v_1'd_1^2 + v_2'd_2^2 + v_3'd_3^2 &= \frac{2}{3} \\ w_1'(c_1b_{11} + c_2b_{12} + c_3b_{13}) + w_2'(c_1b_{21} + c_2b_{22} + c_3b_{23}) + w_3'(c_1b_{31} + c_2b_{32} + c_3b_{33}) + v_1'(\beta_{11}d_1 + \beta_{12}d_2 + \beta_{13}d_3) + v_2'(\beta_{21}d_1 + \beta_{22}d_2 + \beta_{23}d_3) + v_3'(\beta_{31}d_1 + \beta_{32}d_2 + \beta_{33}d_3) &= \frac{2}{3} \end{split}$$

With constraints

$$c_1 = a_{11} + a_{12} + a_{13} = \frac{1}{2}(b_{11} + b_{12} + b_{13})$$

$$c_2 = a_{21} + a_{22} + a_{23} = \frac{1}{2}(b_{21} + b_{22} + b_{23})$$

$$c_3 = a_{31} + a_{32} + a_{33} = \frac{1}{2}(b_{31} + b_{32} + b_{33})$$

$$d_1 = \alpha_{11} + \alpha_{12} + \alpha_{13} = \frac{1}{2}(\beta_{11} + \beta_{12} + \beta_{13})$$

$$d_2 = \alpha_{21} + \alpha_{22} + \alpha_{23} = \frac{1}{2}(\beta_{21} + \beta_{22} + \beta_{23})$$

$$d_3 = \alpha_{31} + \alpha_{32} + \alpha_{33} = \frac{1}{2}(\beta_{31} + \beta_{32} + \beta_{33})$$

This is fourteen (14) equations with fifty-four (54) unknowns. That means the expected scheme it not unique; we can have a family of 3 – stage schemes.

1. Choosing the Parameters

$$\begin{aligned} w_1 &= w_2 = w_3 = 0, \quad v_1 = v_2 = v_3 = \frac{1}{3}, \quad c_1 = d_1 = a_{11} = b_{11} = \alpha_{11} = \beta_{11} = 0 \\ w_1' &= w_2' = \frac{1}{2}, \quad w_3' = 1, \quad v_1' = v_2' = v_3' = 0, \quad c_2 = c_3 = b_{21} = b_{31} = \frac{2}{3} \\ d_2 &= d_3 = a_{12} = b_{12} = \beta_{12} = \beta_{21} = \beta_{31} = \frac{1}{2}, \quad a_{22} = a_{23} = a_{32} = a_{33} = \frac{1}{6}, \\ a_{21} &= a_{31} = b_{22} = b_{23} = b_{32} = b_{33} = \frac{1}{3}, \quad \alpha_{12} = \alpha_{21} = \alpha_{31} = \beta_{22} = \beta_{23} = \beta_{32} = \beta_{33} = \frac{1}{4}, \quad a_{13} = b_{13} = \beta_{13} = -\frac{1}{2}, \quad \alpha_{13} = -\frac{1}{4}, \quad \alpha_{22} = \alpha_{23} = \alpha_{32} = \alpha_{33} = \frac{1}{8} \end{aligned}$$

Then equations (3.3.1) and (3.3.2) becomes

$$y_{n+1} = \frac{y_n + hy_n'}{1 + y_n \frac{1}{3}(H_1 + H_2 + H_3)}$$
(3.3.14)

and

$$y'_{n+1} = y'_n + \frac{1}{h} \left(\frac{1}{2} K_1 + \frac{1}{2} K_2 + K_3 \right)$$
 (3.3.15)

where

$$K_{1} = \frac{h^{2}}{2} f\left(x_{n}, y_{n} + \frac{1}{2}K_{2} - \frac{1}{2}K_{3}, y_{n}' + \frac{1}{h}\left(\frac{1}{2}K_{2} - \frac{1}{2}K_{3}\right)\right)$$

$$K_{2} = K_{3} = \frac{h^{2}}{2} f\left(x_{n} + \frac{2}{3}h, y_{n} + \frac{2}{3}hy_{n}' + \frac{1}{6}(2K_{1} + K_{2} + K_{3}), y_{n}' + \frac{1}{h}\left(\frac{2}{3}K_{1} + \frac{1}{3}K_{2} + \frac{1}{3}K_{3}\right)\right)$$

and

$$H_{1} = \frac{h^{2}}{2} g\left(x_{n}, z_{n} + \frac{1}{4}H_{2} - \frac{1}{4}H_{3}, z'_{n} + \frac{1}{h}\left(\frac{1}{4}H_{2} - \frac{1}{4}H_{3}\right)\right)$$

$$H_{2} = H_{3} = \frac{h^{2}}{2} g\left(x_{n} + \frac{1}{2}h, z_{n} + \frac{1}{2}hz'_{n} + \frac{1}{4}H_{1} + \frac{1}{8}H_{2} + \frac{1}{8}H_{3}, z'_{n} + \frac{1}{h}\left(\frac{1}{2}H_{1} + \frac{1}{4}H_{2} + \frac{1}{4}H_{3}\right)\right)$$

But since $K_2 = K_3$ and $H_2 = H_3$, equations (3.3.14) and (3.3.15) becomes

$$y_{n+1} = \frac{y_n + hy_n'}{1 + y_n \frac{1}{3}(H_1 + 2H_2)}$$
 (3.3.16)

and

$$y'_{n+1} = y'_n + \frac{1}{2h}(K_1 + 3K_2) \tag{3.3.17}$$

where

$$K_{1} = \frac{h^{2}}{2} f(x_{n}, y_{n}, y'_{n}) = \frac{h^{2}}{2} f_{n}$$

$$K_{2} = \frac{h^{2}}{2} f\left(x_{n} + \frac{2}{3}h, y_{n} + \frac{2}{3}hy'_{n} + \frac{1}{3}(K_{1} + K_{2}), y'_{n} + \frac{2}{3h}(K_{1} + K_{2})\right)$$

and

$$H_1 = \frac{h^2}{2} g(x_n, z_n, z_n') = \frac{h^2}{2} g_n$$

$$H_2 = \frac{h^2}{2} g\left(x_n + \frac{1}{2}h, z_n + \frac{1}{2}hz_n' + \frac{1}{4}(H_1 + H_2), z_n' + \frac{1}{2h}(H_1 + H_2)\right)$$

2. Choosing the Parameters

$$\begin{split} w_1 &= w_2 = w_3 = 0, \quad v_1 = \frac{2}{3}, \quad v_2 = v_3 = \frac{1}{6}, \quad c_1 = d_1 = \frac{1}{3}, \quad c_2 = \frac{2-\sqrt{2}}{3}, \quad c_3 = \frac{4+\sqrt{2}}{6}, \\ a_{11} &= a_{22} = a_{33} = b_{11} = \alpha_{11} = \alpha_{22} = \alpha_{33} = \beta_{11} = \beta_{22} = \beta_{33} = \frac{1}{3}, \quad b_{22} = b_{33} = 1 \\ w_1' &= 1, \quad w_2' = \frac{1}{3}, \quad w_3' = \frac{2}{3}, \quad v_1' = v_2' = v_3' = 0, \quad a_{12} = \frac{1}{3}, \quad a_{13} = -\frac{1}{3}, \quad d_2 = \frac{2-\sqrt{6}}{6}, \\ a_{21} &= a_{23} = \frac{1-\sqrt{2}}{6}, \quad a_{31} = a_{32} = \frac{2+\sqrt{2}}{12}, \quad b_{12} = \frac{1+2\sqrt{2}}{9}, \quad b_{13} = \frac{2-2\sqrt{2}}{9}, \quad b_{23} = \frac{-2\sqrt{2}}{3}, \\ b_{21} &= b_{31} = \frac{1}{3}, \quad b_{32} = \frac{\sqrt{2}}{3}, \quad \alpha_{12} = \frac{1}{3}, \quad \alpha_{13} = -\frac{1}{3}, \quad \alpha_{21} = \alpha_{23} = -\frac{\sqrt{6}}{12}, \quad \alpha_{31} = \frac{\sqrt{6}}{12}, \\ \alpha_{32} &= \frac{\sqrt{6}}{6}, \quad \beta_{12} = \beta_{13} = \frac{1}{6}, \quad \beta_{21} = \beta_{23} = \frac{1-\sqrt{6}}{6}, \quad \beta_{31} = \beta_{32} = \frac{1+\sqrt{6}}{6}, \quad d_3 = \frac{2+\sqrt{6}}{6}, \end{split}$$

Then equations (3.3.1) and (3.3.2) becomes

$$y_{n+1} = \frac{y_n + hy_n'}{1 + \frac{1}{6}y_n(4H_1 + H_2 + H_3)}$$
(3.3.18)

and

$$y'_{n+1} = y'_n + \frac{1}{3h}(3K_1 + K_2 + 2K_3)$$
(3.3.19)

where

$$\begin{split} K_1 &= \frac{h^2}{2} f \left(x_n + \frac{1}{3} h, y_n + \frac{1}{3} h y_n' + \frac{1}{3} (K_1 + K_2 - K_3), y_n' + \frac{1}{3h} \left(K_1 + \left(\frac{1 + 2\sqrt{2}}{3} \right) K_2 + \left(\frac{2 - \sqrt{2}}{3} \right) K_3 \right) \right) \\ K_2 &= \frac{h^2}{2} f \left(x_n + \left(\frac{2 - \sqrt{2}}{3} \right) h, y_n + \left(\frac{2 - \sqrt{2}}{3} \right) h y_n' + \left(\frac{1 - \sqrt{2}}{6} \right) K_1 + \frac{1}{3} K_2 + \left(\frac{1 - \sqrt{2}}{6} \right) K_3, y_n' + \frac{1}{h} \left(\frac{1}{3} K_1 + K_2 - \frac{2\sqrt{2}}{2} K_3 \right) \right) \\ K_3 &= \frac{h^2}{2} f \left(x_n + \left(\frac{4 + \sqrt{2}}{6} \right) h, y_n + \left(\frac{4 + \sqrt{2}}{6} \right) h y_n' + \left(\frac{2 + \sqrt{2}}{12} \right) K_1 + \left(\frac{2 + \sqrt{2}}{12} \right) K_2 + \frac{1}{3} K_3, y_n' + \frac{1}{h} \left(\frac{1}{3} K_1 + \frac{2\sqrt{2}}{2} K_2 + K_3 \right) \right) \end{split}$$

and

$$H_{1} = \frac{h^{2}}{2}g\left(x_{n} + \frac{1}{3}h, z_{n} + \frac{1}{3}hz'_{n} + \frac{1}{3}(H_{1} + H_{2} + H_{3}), z'_{n} + \frac{1}{6h}(2H_{1} + H_{2} + H_{3})\right)$$

$$H_{2} = \frac{h^{2}}{2}g\left(x_{n} + \left(\frac{2 - \sqrt{6}}{6}\right)h, z_{n} + \left(\frac{2 - \sqrt{6}}{6}\right)hz'_{n} - \frac{1}{6}(\sqrt{6}H_{1} - 2H_{2} + \sqrt{6}H_{3}), z'_{n}\right)$$

$$+ \frac{1}{6h}\left(\left(\sqrt{1 - 6}\right)H_{1} + 2H_{2} + \left(\sqrt{1 - 6}\right)H_{3}\right)\right)$$

$$H_{3} = \frac{h^{2}}{2}g\left(x_{n} + \left(\frac{2 + \sqrt{6}}{6}\right)h, z_{n} + \left(\frac{2 + \sqrt{6}}{6}\right)hz'_{n} + \frac{1}{6}\left(\sqrt{6}H_{1} + \sqrt{6}H_{2} + 2H_{3}\right), z'_{n}\right)$$

$$+ \frac{1}{6h}\left(\left(\sqrt{1 + 6}\right)H_{1} + \left(\sqrt{1 + 6}\right)H_{2} + 2H_{3}\right)\right)$$

CHAPTER FOUR

ANALYSIS OF THE METHODS

4.0: Introduction

There are two types of error involved in a Runge–Kutta method: *roundoff* error and *truncation* error (also known as *discretization* error). Round-off error occurs usually when the method is implemented on a computer. Normally, round-off error is not considered in the numerical analysis of the algorithm, since it depends on the computer on which the algorithm is implemented, and thus is external to the numerical algorithm, Julyan and Oreste (1992). Truncation error is caused by truncation of the infinite Taylor series during the development of the new formula.

An obvious requirement for a successful numerical scheme is that the truncation error should be as small as is desired by using a sufficiently small step length: this concept is known as *convergence*. A concept closely related to convergence is known as *consistency*; a method is said to be consistent (with the initial value problem) if

$$\phi(x_n, y_n, y_n'; 0) = f(x_n, y_n, y_n') \tag{4.1}$$

Where $\emptyset(x_n, y_n, y'_n; h)$ is defined by (3.1)

The two crucial concepts in the analysis of numerical error are *local* truncation error and global truncation (discretization) error. The Truncation error is defined by

$$T_{n+1} = y(x_{n+1}) - y_{n+1} = |y_{n+1} - y_n - h\emptyset(x_n, y_n, y_n'; h)|$$
(4.2)

Where $y(x_{n+1})$ is the exact equation of the differential equation (1.1) and y_{n+1} is the approximate solution

4.1: THE TRUNCATION ERROR OF THE METHOD

Thus, the Truncation error of the method is given by

$$T_{n+1} = y(x_{n+1}) - \frac{y_n + hy_n + \sum_{r=1}^{s} w_r k_r}{1 + y_n \sum_{r=1}^{s} v_r H_r}$$
(4.1.1)

Using the Taylor's series expansion of $y(x_{n+1})$, K_r , and H_r about (x_n, y_n, y'_n) . The local truncation error of one – stage scheme (3.1.32) of order three is given as

$$T_{n+1} = \left[\left(\frac{1}{24} - \frac{c_1^2}{4} \right) \Delta^2 f_n + \left(\frac{1}{24} - \frac{c_1}{4} \right) \left(2\Delta f_n f_{y'} + f_n y \right) \right] h^4 \tag{4.1.2}$$

If $c_1 = 0$, then

$$T_{n+1} = \left[\Delta^2 f_n + 2\Delta f_n f_{y'} + f_n f_y\right] h^4 \tag{4.1.3}$$

Where

$$\Delta^2 f_n = f_{xx} + y'_n^2 f_{yy} + f^2 f_{y'y'} + 2y' f_n f_{yy'} + 2f_n f_{xy'}, \text{ and}$$

$$\Delta f_n = f_x + y' f_y + f_n f_{y'}$$

By Lotkin (1957) technique, we can find a bound for the partial derivative of f_n from the inequality defined by

$$f_n < M$$
 and $\left| \frac{\partial^{i+j+k}}{\partial x^i \partial y^j \partial y'^k} f(x_n, y_n, y'_n) \right| < \frac{L^{i+j+k}}{M^{j+k-1}}$ (4.1.4)

Where M and L are constant and $i + j + k \le M$. This condition is used by Lotkin (1957). Thus using (3.4.6) we have

$$f_x < M$$
, $f_y < L$, $f_{y\prime} < L$, $f_{xx} < L^2 M$, $f_{xy} < L^2$, $f_{xy\prime} < L^2$
$$f_{yy} < \frac{L^2}{M}$$
, $f_{yy\prime} < \frac{L^2}{M}$

So, the error bound for (3.4.5) will be

$$T_{n+1} < \left[\left| \frac{1}{24} \right| |\Delta^2 f_n| + |2| \left| \Delta f_n f_{y'} \right| + \left| f_n f_y \right| \right] h^4 < \frac{1}{24} (10L^2 M + LM) h^4$$
(4.1.5)

Hence the local truncation error is

$$T_{n+1} < \frac{1}{24} (10L^2M + LM)h^4$$

To estimate the step size, we can keep

$$T_{n+1} < \frac{1}{24} (10L^2M + LM)h^4 < TOL$$
 (4.1.6)

Where *TOL* is an allowance tolerance to be specified by the user.

4.2: CONVERGENCE

A numerical method is said to be convergent if the numerical solution approaches the exact solution as the step size tends to zero.

Convergent =
$$\lim_{n\to 0} |y(x_{n+1}) - y_{n+1}|$$

In other words, if the *discretiation* error at x_{n+1} tends to zero as $h \to \infty$, i.e if

$$e_{n+1} = |y(x_{n+1}) - y_{n+1}| \to 0 \text{ as } n \to \infty$$
 (4.2.1)

From equation (3.1.33),

$$y'_{n+1} = \frac{y'_n}{1 + \frac{2}{h} y'_n H_1} \tag{4.2.2}$$

And the numerical solution $y'(x_{n+1})$ seems to satisfy the equation

$$y'(x_{n+1}) = \frac{y'_n}{1 + \frac{2}{h}y'(x_n)H_1} + T_{n+1}$$
(4.2.3)

Where T_{n+1} is a local truncation error.

Subtracting equation (4.2.3) from (4.2.2) gives

$$y'_{n+1} - y'(x_{n+1}) = \frac{y'_n}{1 + \frac{2}{h}y'_n H_1} - \frac{y'_n}{1 + \frac{2}{h}y'(x_n)H_1} + T_{n+1}$$
(4.2.4)

Adopting equation (4.2.4) gives

$$y'_{n+1} - y'(x_{n+1}) = \frac{y'_n(1 + \frac{2}{h}y'_n H_1) - y'_n(1 + \frac{2}{h}y'(x_n) H_1)}{(1 + \frac{2}{h}y'_n H_1)(1 + \frac{2}{h}y'(x_n) H_1)} + T_{n+1}$$
(4.2.5)

Expanding the brackets and re-arranging gives

$$y'_{n+1} - y'(x_{n+1}) = \frac{\frac{2}{h} y'_n H_1(y'_n - y'(x_n))}{\left(1 + \frac{2}{h} y'_n H_1\right) \left(1 + \frac{2}{h} y'(x_n) H_1\right)} + T_{n+1}$$

This implies that

$$e_{n+1} = \frac{e(\frac{2}{h}y'_n H_1)}{(1 + \frac{2}{h}y'_n H_1)(1 + \frac{2}{h}y'(x_n) H_1)} + T_{n+1}$$
(4.2.6)

From equations (4.2.6), setting

$$A_n = \left(\frac{2}{h}y_n'H_1\right), B_n = \left(1 + \frac{2}{h}y_n'H_1\right), C_n = \left(1 + \frac{2}{h}y'(x_n)H_1\right) \text{ and } T_{n+1} = T$$

Then

$$e_{n+1} = \frac{A_n}{B_n C_n} e_n + T \tag{4.2.7}$$

Let $B = \max B_n > 0$, $C = \max C_n > 0$ and $A = \max A_n < 0$ then (4.2.7) becomes,

$$e_{n+1} \le \frac{A}{BC}e_n + T$$

Set $\frac{A}{BC} = K < 1$, then

$$e_{n+1} \le Ke_n + T \tag{4.2.8}$$

If n = 0, then from (4.2.8),

$$e_1 = Ke_0 + T$$

 $e_2 = Ke_1 + T = K^2e_0 + KT + T$ by substituting the value of e_1

$$e_3 = Ke_2 + T = K^3e_0 + K^2T + T$$

Continuing in this manner, we get the following

$$e_{n+1} = K^{n+1}e_0 + \sum_{t=0}^{n+1} K^t T \tag{4.2.9}$$

Since $\frac{A}{BC} = K < 1$, then one can see that as $n \to \infty$, $e_{n+1} \to 0$. This proves that the scheme converges.

4.3: CONSISTENCY

A scheme is said to be consistent, if the difference equation of the computation formula exactly approximates the differential equation it intends to solve as the step size ends to zero. To prove if equation (3.1.32) is consistent, subtract y_n from both side of (3.1.32), then

$$y_{n+1} - y_n = \frac{y_n + hy_n'}{1 + y_n H_1} - y_n \tag{4.3.1}$$

Therefore

$$y_{n+1} - y_n = \frac{y_n + hy_n' - y_n(1 + y_n H_1)}{1 + y_n H_1}$$

Simplifying gives

$$y_{n+1} - y_n = \frac{hy_n' - y_n^2 H_1}{1 + y_n H_1} \tag{4.3.2}$$

But since

$$H_1 = \frac{h^2}{2} g\left(x_n + \frac{1}{3}h, z_n + \frac{1}{3}hz'_n + \frac{1}{3}H_1, z'_n + \frac{2}{3h}H_1\right)$$

Then equation (4.3.2) becomes

$$y_{n+1} - y_n = \frac{hy_n' - y_n^2 \left[\frac{h^2}{2} g\left(x_n + \frac{1}{3} h, z_n + \frac{1}{3} h z_n' + \frac{1}{3} H_1, z_n' + \frac{2}{3h} H_1 \right) \right]}{1 + y_n \left[\frac{h^2}{2} g\left(x_n + \frac{1}{3} h, z_n + \frac{1}{3} h z_n' + \frac{1}{3} H_1, z_n' + \frac{2}{3h} H_1 \right) \right]}$$

Dividing all through by h and taking the limit as h tends to zero on both sides gives

$$\lim_{h \to \infty} \frac{y_{n+1} - y_n}{h} = y_n'$$

Hence, the method is consistent.

CHAPTER FIVE

NUMERICAL EXAMPLES

5.0 INTRODUCTION

Having gotten our schemes, part of the objectives is to apply the scheme derived to approximate real life problems. Second order linear differential equations have a variety of applications in science and engineering. In this research explore one of them, the spring problems.

5.1 Numerical Problems

Example 1.

A spring of mass 2kg has a damping constant of 8, and a force of 5N is required to keep the spring stretched 0.5m beyond its natural length. The spring is stretched 1m beyond its natural length and then release with zero velocity. Find the position of mass at time t.

Solution

The problem above can be modeled by the initial value problem

$$\frac{d^2x}{dt^2} + 4\frac{dx}{dt} + 5x = 0, \quad x(0) = 1, \quad x'(0) = 0$$

The exact solution is

$$x(t) = e^{-2t} \left(2\sin t + \cos t \right)$$

Example 2.

A spring of mass 2kg has a damping constant of 10, and a force of 6N is required to keep the spring stretched 0.5m beyond its natural length. The spring

is stretched 1m beyond its natural length and then release with zero velocity. Find the position of mass at time t.

Solution

The above spring problem is modeled by the initial value problem

$$\frac{d^2x}{dt^2} + 5\frac{dx}{dt} + 6x = 0, \quad x(0) = 1, \quad x'(0) = 0$$

The exact solution is

$$x(t) = -2e^{-3t} + 3e^{-2t}$$

Example 3. (Variable Coefficient Equation)

Consider the equation $y'' = (1 + x^2)y$, y(0) = 1, y'(0) = 0, $x \in [0,1]$

The exact solution is

$$y(x) = e^{x^2/2}$$

[Sources: Jain 1984]

Example 4. (Non-linear Differential Equation)

Consider a non-linear ordinary differential equation

$$y'' - x(y')^2 = 0$$
, $y(0) = 1$, $y'(0) = \frac{1}{2}$

The exact solution is

$$y(x) = 1 + \frac{1}{2} \ln \left(\frac{2+x}{2-x} \right)$$

[Sources: Jacob (2010), Taiwo and Osilagun (2011) and Anake et-al (2012)]

5.2 PRESENTATION OF THE RESULTS

The results below are the exact and the approximate solution to each problem.

(Note that. All results is in 10 decimal places)

Table 5.2.1 results of example 1 at h = 0.1

	Exact Solutions		Approxima	Approximate Solutions		ors
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}'_n	\mathcal{Y}_n	y'_n
0.1	0.978113886	-0.408683442	0.971852523	-0.406934307	6.26E-03	-1.75E-03
0.2	0.923302344	-0.665860175	0.901166794	-0.661027045	2.21E-02	-4.83E-03
0.3	0.848669638	-0.810924641	0.801778305	-0.798072019	4.69E-02	1.29E-02
0.4	0.763813263	-0.874884702	0.682786227	-0.847056781	8.10E-02	-2.78E-02
0.5	0.675586181	-0.881853996	0.548234367	-0.830600061	1.27E-01	-5.13E-02
0.6	0.588720400	-0.850335224	0.393287231	-0.765107383	1.95E-01	-8.52E-02
0.7	0.506332013	-0.794310629	0.172338281	-0.659316519	3.34E-01	-1.35E-01
0.8	0.430326053	-0.724158485	0.636264391	-0.500370482	-2.06E-01	-2.24E-01
0.9	0.361717571	-0.647415336	0.535303951	-0.585436941	-1.74E-01	-6.20E-02
1.0	0.300883394	-0.569403570	0.406425080	-0.599863195	-1.06E-01	3.05E-02

Table 5.2.2 results of example 1 at h = 0.01

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}'_n	\mathcal{Y}_n	y'_n
0.01	0.999753311	-0.049009117	0.999746676	-0.049007072	6.63E-06	-2.05E-06
0.02	0.999026303	-0.096072539	0.999000136	-0.096068285	2.62E-05	-4.25E-06
0.03	0.997838168	-0.141243491	0.997779902	-0.141236237	5.83E-05	-7.25E-06
0.04	0.996207568	-0.18457404	0.996104951	-0.184562407	1.03E-04	-1.16E-05
0.05	0.994152652	-0.226115112	0.993993728	-0.226097174	1.59E-04	-1.79E-05
0.06	0.991691064	-0.265916514	0.991464159	-0.265889831	2.27E-04	-1.79E-05
0.07	0.988839958	-0.304026952	0.988533663	-0.303988608	3.06E-04	-3.83E-05
0.08	0.985616003	-0.340494051	0.985219165	-0.340440684	3.97E-04	-5.34E-05
0.09	0.982035399	-0.375364374	0.981537104	-0.375292206	4.98E-04	-7.22E-05
0.1	0.978113886	-0.408683442	0.977503449	-0.408588307	6.10E-04	-9.51E-05

Table 5.2.3 results of example 1 at h = 0.001

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}_n'	\mathcal{Y}_n	y'_n
0.001	0.999997503	-0.004990009	0.999997497	-0.004990007	6.66E-09	-2.08E-09
0.002	0.999990027	-0.009960073	0.999990000	-0.009960069	2.66E-08	-4.18E-09
0.003	0.999977590	-0.014910247	0.999977530	-0.014910241	5.98E-08	-6.35E-09
0.004	0.999960213	-0.019840585	0.999960107	-0.019840577	1.06E-07	-8.68E-09
0.005	0.999937915	-0.024751143	0.999937749	-0.024751131	1.66E-07	-1.12E-08
0.006	0.999910717	-0.029641974	0.999910478	-0.029641959	2.39E-07	-1.40E-08
0.007	0.999878638	-0.034513132	0.999878313	-0.034513115	3.24E-07	-1.72E-08
0.008	0.999841697	-0.039364673	0.999841274	-0.039364652	4.23E-07	-2.08E-08
0.009	0.999799915	-0.044196650	0.999799380	-0.044196625	5.35E-07	-2.48E-08
0.01	0.999753311	-0.049009117	0.999752650	-0.049009087	6.60E-07	-2.94E-08

Table 5.2.4 results of example 2 at h = 0.1

	Exact S	Exact Solutions		te Solutions	Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}'_n	y_n	y_n'
0.1	0.974555818	-0.467475194	0.965148944	-0.490368272	9.41E-03	2.29E-02
0.2	0.913336866	-0.729050460	0.878800789	-0.782015951	3.45E-02	5.30E-02
0.3	0.833295589	-0.853451858	0.762787276	-0.923295925	7.05E-02	6.98E-02
0.4	0.745598469	-0.888808513	0.632219564	-0.955356931	1.13E-01	6.65E-02
0.5	0.657378003	-0.868495686	0.494344330	-0.911516387	1.63E-01	4.30E-02
0.6	0.572984859	-0.815371942	0.342244979	0.816304591	2.31E-01	9.33E-04
0.7	0.494878035	-0.744843214	-1.45945457	-0.681774224	1.95E+00	-6.31E-02
0.8	0.424253647	-0.667071388	-1.51700728	0.286422849	1.94E+00	-9.53E-01
0.9	0.361485639	-0.588560253	-1.44948531	0.924225035	1.81E+00	-1.51E+00
1.0	0.306431713	-0.513289289	-1.30394335	1.292676829	1.61E+00	-1.81E+00

Table 5.2.5 results of example 2 at h = 0.01

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}'_n	\mathcal{Y}_n	y'_n
0.01	0.999704953	-0.058518839	0.999694999	-0.058810937	9.95E-06	2.92E-04
0.02	0.998839250	-0.114149433	0.998797086	-0.115138120	4.22E-05	9.89E-04
0.03	0.997431230	-0.167000090	0.997331275	-0.169050819	1.00E-04	2.05E-03
0.04	0.995508166	-0.217175458	0.995321915	-0.220616869	1.86E-04	3.44E-03
0.05	0.993096301	-0.264776650	0.992792696	-0.269902694	3.04E-04	5.13E-03
0.06	0.990220887	-0.309901352	0.989766661	-0.316973325	4.54E-04	7.07E-03
0.07	0.986906214	-0.352643937	0.986266216	-0.361892421	6.40E-04	9.25E-03
0.08	0.983175645	-0.393095567	0.982313139	-0.404722292	8.63E-04	1.16E-02
0.09	0.979051646	-0.431344302	0.977928587	-0.445523915	1.12E-03	1.42E-02
0.1	0.974555818	-0.467475194	0.973133113	-0.484356955	1.42E-03	1.69E-02

Table 5.2.6 results of example 2 at h = 0.001

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}_n'	\mathcal{Y}_n	y'_n
0.001	0.999997005	-0.005985019	0.999996995	-0.005988011	1.00E-08	2.99E-06
0.002	0.999988040	-0.011940152	0.999987997	-0.011950589	4.29E-08	1.04E-05
0.003	0.999973135	-0.017865512	0.999973031	-0.017887807	1.03E-07	2.23E-05
0.004	0.999952319	-0.023761212	0.999952124	-0.023799736	1.95E-07	3.85E-05
0.005	0.999925622	-0.029627365	0.999925299	-0.029686448	3.23E-07	5.91E-05
0.006	0.999893074	-0.035464083	0.999892583	-0.035548015	4.91E-07	8.39E-05
0.007	0.999854704	-0.041271478	0.999854000	-0.041384508	7.03E-07	1.13E-04
0.008	0.999810541	-0.047049662	0.999809576	-0.047195999	9.64E-07	1.46E-04
0.009	0.999760614	-0.052798745	0.999759336	-0.052982558	1.28E-06	1.84E-04
0.01	0.999704953	-0.058518839	0.999703304	-0.058744256	1.65E-06	2.25E-04

Table 5.2.7 results of example 3 at h = 0.1

	Exact S	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}_n'	\mathcal{Y}_n	\mathcal{Y}_n'	
0.1	1.005012521	0.100501252	1.005025126	0.100250417	-1.26E-05	2.51E-04	
0.2	1.020201340	0.204040268	1.020129605	0.201733389	7.17E-05	2.31E-03	
0.3	1.046027860	0.313808358	1.045542449	0.306464227	4.85E-04	7.34E-03	
0.4	1.083287068	0.433314827	1.081707487	0.416660651	1.58E-03	1.67E-02	
0.5	1.133148453	0.566574227	1.129290062	0.534778113	3.86E-03	3.18E-02	
0.6	1.197217363	0.718330418	1.189200895	0.663611165	8.02E-03	5.47E-02	
0.7	1.277621313	0.894334919	1.262631449	0.806411362	1.50E-02	8.79E-02	
0.8	1.377127764	1.101702211	1.351102688	0.967028171	2.60E-02	1.35E-01	
0.9	1.499302500	1.349372250	1.456529654	1.150081082	4.28E-02	1.99E-01	
1.0	1.648721271	1.648721271	1.581305030	1.361173403	6.74E-02	2.88E-01	

Table 5.2.8 results of example 3 at h = 0.01

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}'_n	\mathcal{Y}_n	y'_n
0.01	1.000050001	0.010000500	1.000050003	0.01000025	-1.25E-09	2.50E-07
0.02	1.000200020	0.020004000	1.000200013	0.020001722	7.08E-09	2.28E-06
0.03	1.000450101	0.030013503	1.000450054	0.030013503	4.74E-08	7.14E-06
0.04	1.000800320	0.040032013	1.000800169	0.040016172	1.51E-07	1.58E-05
0.05	1.001250782	0.050062539	1.001250421	0.050033155	3.61E-07	2.94E-05
0.06	1.001801621	0.060108097	1.001800894	0.060059320	7.27E-07	4.88E-05
0.07	1.002453004	0.070171710	1.002451691	0.070096677	1.31E-06	7.50E-05
0.08	1.003205125	0.080256410	1.003202935	0.080147243	2.19E-06	1.09E-04
0.09	1.004058212	0.090365239	1.004054772	0.090213038	3.44E-06	1.52E-04
0.1	1.005012521	0.100501252	1.005007363	0.100296091	5.16E-06	2.05E-04

Table 5.2.9 results of example 3 at h = 0.001

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}_n'	\mathcal{Y}_n	y'_n
0.001	1.000000500	0.001000001	1.000000500	0.001000000	0.00E-00	0.00E-00
0.002	1.000002000	0.002000004	1.000002000	0.002000002	0.00E-00	2.28E-09
0.003	1.000004500	0.003000014	1.000004500	0.003000006	0.00E-00	7.14E-09
0.004	1.000008000	0.004000032	1.000008000	0.004000016	0.00E-00	1.58E-08
0.005	1.000012500	0.005000063	1.000012500	0.005000033	0.00E-00	2.94E-08
0.006	1.000018000	0.006000108	1.000018000	0.006000059	0.00E-00	4.87E-08
0.007	1.000024500	0.007000172	1.000024500	0.007000097	0.00E-00	7.49E-08
0.008	1.000032001	0.008000256	1.000032000	0.008000147	0.00E-00	1.09E-07
0.009	1.000040501	0.009000365	1.000040500	0.009000213	0.00E-00	1.52E-07
0.01	1.000050001	0.010000500	1.000050001	0.010000296	0.00E-00	2.05E-07

Table 5.2.10 results of example 4 at h = 0.1

	Exact S	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}_n'	\mathcal{Y}_n	\mathcal{Y}'_n	
0.1	1.050041729	0.501253133	1.050073940	0.503168650	-3.22E-05	-1.92E-03	
0.2	1.100335348	0.505050505	1.100577720	0.509928057	-2.42E-04	-4.88E-03	
0.3	1.151140436	0.511508951	1.151877265	0.520516997	-7.37E-04	-9.01E-03	
0.4	1.202732554	0.520833333	1.204368009	0.535350750	-1.64E-03	-1.45E-02	
0.5	1.255412812	0.533333333	1.258494045	0.555072830	-3.08E-03	-2.17E-02	
0.6	1.309519604	0.549450549	1.314773677	0.580648013	-5.25E-03	-3.12E-02	
0.7	1.365443754	0.569800570	1.373836442	0.613531812	-8.39E-03	-4.37E-02	
0.8	1.423648930	0.595238095	1.436481545	0.656001612	-1.28E-02	-6.08E-02	
0.9	1.484700279	0.626959248	1.503780137	0.711893892	-1.91E-02	-8.49E-02	
1.0	1.549306144	0.666666667	1.577285372	0.788697022	-2.80E-02	-1.22E-01	

Table 5.2.11 results of example 4 at h = 0.01

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}'_n	\mathcal{Y}_n	y'_n
0.01	1.005000042	0.500012500	1.005000057	0.500025667	-1.57E-08	-1.32E-05
0.02	1.010000333	0.500050005	1.010000457	0.500077340	-1.24E-07	-2.73E-05
0.03	1.015001125	0.500112525	1.015001460	0.500155033	-3.35E-07	-4.25E-05
0.04	1.020002667	0.500200080	1.020003326	0.500258771	-6.59E-07	-5.87E-05
0.05	1.025005210	0.500312695	1.025006315	0.500388585	-1.10E-06	-7.59E-05
0.06	1.030009005	0.500450405	1.030010689	0.500544515	-1.68E-06	-9.41E-05
0.07	1.035014302	0.500613251	1.035016709	0.500726612	-2.41E-06	-1.13E-04
0.08	1.040021354	0.500801282	1.040024636	0.500801282	-3.28E-06	-1.34E-04
0.09	1.045030412	0.501014554	1.045034733	0.501169545	-4.32E-06	-1.55E-04
0.1	1.050041729	0.501253133	1.050047264	0.501430525	-5.53E-06	-1.77E-04

Table 5.2.12 results of example 4 at h = 0.001

	Exact Solutions		Approximate Solutions		Errors	
t	$y(x_n)$	$y'(x_n)$	\mathcal{Y}_n	\mathcal{Y}_n'	\mathcal{Y}_n	y_n'
0.001	1.000500000	0.500000125	1.000500000	0.500000251	0.00E-00	-1.26E-07
0.002	1.001000000	0.500000500	1.001000000	0.500000752	0.00E-00	-2.52E-07
0.003	1.001500001	0.500001125	1.001500001	0.500001505	0.00E-00	-3.80E-07
0.004	1.002000003	0.500002000	1.002000003	0.500002509	0.00E-00	-5.09E-07
0.005	1.002500005	0.500003125	1.002500006	0.500003763	0.00E-00	-6.38E-07
0.006	1.003000009	0.500004500	1.003000010	0.500005269	-1.37E-09	-7.69E-07
0.007	1.003500014	0.500006125	1.003500016	0.500007026	-1.90E-09	-9.01E-07
0.008	1.004000021	0.500008000	1.004000024	0.500009033	-2.53E-09	-1.03E-06
0.009	1.004500030	0.500010125	1.004500034	0.500011292	-3.24E-09	-1.17E-06
0.01	1.005000042	0.500012500	1.005000046	0.500013802	-4.05E-09	-1.30E-06

Table 5.2.13 Comparing errors of example 4, Jacob (2010) and new scheme at $h = 3.125 \times 10^{-3}$

	Exact Solutions	Computed with new scheme	Absolute Errors in Jacob (2010)	Absolute error in new Scheme
t	$y(x_n)$	y_n	y_n	\mathcal{Y}_n
0.1	1.001562501	1.001562502	6.125E-08	4.41E-10
0.2	1.003125010	1.003125014	1.211E-07	3.51E-09
0.3	1.004687534	1.004687544	1.874E-07	9.31E-09
0.4	1.006250081	1.006250099	2.616E-07	1.79E-08
0.5	1.007812659	1.007812688	3.534E-07	2.95E-08

5.3 Discussions

Example 1 and 2 are second order linear ordinary differential equations, the schemes favorably approximate the exact solution of the differential equations. The results are highlighted in tables (5.2.1) through (5.2.6) with various step lengths. Thus, the results are better with smaller step length.

In table (5.2.7), (5.2.8) and (5.2.9) are results obtained by applying the schemes to example 3, the variable coefficient differential equation. The result performed well and approximate the exact solution better as the step size goes to h = 0.001.

Example 4 is a non-linear ordinary differential equation which is also well approximated with results in tables (5.2.10), (5.2.11) and (5.2.12). The results show superiority over the results of Optimal Order Method of Jacob (2010).

CHAPTER SIX

SUMMARY, RECOMMENDATION AND CONCLUSION

6.0 Summary

The 3 – State Implicit Rational Runge-Kutta Scheme was derived and used for the solution of problems leading to second order ordinary differential equation and was found out to be effective. The scheme was found to be consistent, convergent and give better approximation.

6.1 Conclusion

The new numerical schemes derived follows the techniques of rational form of Runge – Kutta methods proposed by Hong (1982) which was adopted by Okunbor (1987) and Ademiluyi and Babatola (2000) by using Taylor and Binomial expansion in stages evaluation. The order condition obtained in this research is up to five (5) and the stage is up to three (3). This is an improvement on the work of earlier authors.

The new schemes are of high accuracy for direct numerical solution of general second order ordinary differential equations. The steps to the derivation of the new schemes are presented in the methodology while the analysis of the schemes proved to be consistent, convergent, the results proves to be good estimate of the exact equations. Thus, the scheme is effective and efficient, these suggest a wider application of the schemes for even more complicated physical problems; since the methods is used to solve equations of the form y'' = (x, y, y') and y'' = (x, y) favorably. Equations of constant and variable coefficients are also considered and the fact that is also used to solve non-linear problem.

6.2 Recommendation

Since the scheme is capable of approximating various problems, we recommend further research in the area to explore and derive more schemes from the rational function of Runge – Kutta methods defined in (3.1) for solving second order ordinary differential equations as in the case of first order equations where many authors worked in that direction.

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